

CHM 335 - Physical Chemistry Laboratory
Course Syllabus
Fall 2017

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3 Text Book

Experiments in Physical Chemistry, Garland, Nibler and Shoemaker (McGraw-Hill, 2017)
ISBN-10:130732307, URI Package, Available only in the URI Bookstore

4 Classroom

280 Beaupre

5 Laboratory Safety

You must wear safety glasses with side guards or safety goggles and have them on before you enter the laboratory. Of course, you must continue to wear them while in the lab. You must also wear a laboratory coat, and you may not wear open-toed shoes or sandals. If you are ever in the laboratory without eye protection or a lab coat, you will have 5 points deducted from your next laboratory report. Repeated offenses will result in failure in the course. Be sure to bring your eye protection and lab coat to the second class meeting, so that you will be allowed to perform the first experiment of the semester.

All students must complete the Department Safety and Environmental Compliance Form. All students must have their Medical Information Form with them during every laboratory period.

6 Course Requirements

This semester you are to perform nine experiments. The nine experiments are divided into two groups of four plus a final experiment at the end of the semester. The first four experiments are designed to gradually teach each of you how to extract data and work up an analysis as a laboratory report. Each of the first four experiments are performed by every member of a section in the same laboratory period.

Some of the experiments are to be performed with a laboratory partner. **Even though you may collect data with a partner, you are responsible for your own laboratory reports. Identical or nearly identical laboratory reports from two or more students are treated as cases of plagiarism.**

In the first experiment you will study the spectrum of a conjugated dye. Again you will work with partners. You will extract data by measuring the spectra and comparing the location of the wavelength maximum with theoretical predictions.

The second experiment is a study of the diffusion of ionic species in water. This experiment is performed with a partner. In this experiment you extract data and confirm that concentration gradients migrate with a particular functional dependence on the time. Details are given in section 13.2.

In the third experiment you measure the viscosity of polyvinyl alcohol in water. The purpose of this experiment is to instruct you in the principles of experimental error. You repeat your measurements many times to observe the statistical distribution of error. Your laboratory reports on this experiment emphasize the analysis of these errors. Details are given in Section 14.2.

In the final of the first four experiments, you again measure the viscosity of polyvinyl alcohol solutions, but after chemical degradation. The purpose is to learn structural information about the polymer from bulk measurements. The laboratory report from this experiment incorporates all sections learned in the first three experiments, and constitutes a complete laboratory report.

The second set of four experiments are performed with partners. Each of the four experiments is performed in every laboratory period by one set of partners. The equipment is rotated among groups during this five week period. The experiments are 1) the vapor pressure of a liquid; 2) the surface tension of butanol solutions; 3) the phase transitions in ammonium chloride; and 4) the heat capacity ratio of gases. Complete laboratory reports are to be submitted for each of these experiments.

In the ninth and final experiment of the semester you study the the partial molar volume of aqueous sodium chloride solutions. The laboratory report for this experiment is written in class, and the report can be viewed as a kind of final exam where the questions are known ahead of time. All materials including paper and graph paper are provided for you.

7 Data Collection

Where possible, all data generated in the laboratory should be entered into a laboratory notebook. Laboratory notebooks can be purchased in the bookstore, and the bookstore notebooks generate carbonless copies, so that a copy of your data is automatically generated. In some instances the data you generate cannot be entered into the notebook. Your TA will inform you when the notebooks are not required. The copy of your data signed by your TA should be included in the data portions of your laboratory reports.

8 Grading

For the first 8 experiments, each of the laboratory report are worth 100 points. These reports, **which are to be submitted both as a hard copy and uploaded electronically**, are due exactly one week after the completion of a given experiment. The electronic versions should be uploaded before coming to class. If, because of a holiday, classes do not meet one week following the completion of a report, the report is due the day following the holiday or on the date the report would normally be due. For example, this semester Wednesday, November 15 follows a Monday schedule. The reports for any experiment performed on Friday, November 8 are due on Wednesday, November 15 even though students in Wednesday's section do not meet again until Wednesday, November 22. In a similar fashion, if a class is canceled because of weather (or any other reason) the report is due the day following the canceled class or the next date the University is officially open. **Reports are due at the beginning of the class on the due date. Late reports will be accepted for grading up to one week after the due date. However, late reports are assessed a penalty of 2 points for each day that they are late. Reports submitted more than one week past the**

due date will not be accepted for grading, and such reports will receive a score of 0. Your laboratory reports must show all calculations, not just sample calculations of each type. The final laboratory report, written in class, is worth 200 points.

Your grades are determined on an absolute scale. The scale is

A : 930 to 1000 points

A- : 900 to 929

B+ : 870 to 899

B : 830 to 869

B- : 800 to 829

C+ : 770 to 799

C : 730 to 769

C- : 700 to 729

D+ : 670 to 699

D : 600 to 669

F : < 600

9 Outline of the Laboratory Procedure

Before being allowed to begin any of the experiments in this course, you are required to prepare an outline of the experimental procedure. You should prepare the outline at home, and bring it to class with you. Your TA then signs and dates the outline before you begin the experiment. The signed outline must be attached to your laboratory report. No laboratory reports can be graded without the attached signed outline.

At the completion of an experiment, you must present your data to your TA. The TA then signs and dates the data. The original signed data must be attached to your laboratory reports for grading. No laboratory report can be graded without the signed data attached.

10 Plagiarism

One of our goals in this course is to reinforce the importance of scientific integrity. In recent years, there have been numerous examples of established scientists generating falsified data or copying material from another source. Acts of plagiarism both damage science and can have important impacts on society. The possibly falsified data associated with the connection between childhood vaccines and autism is an important recent example that has adversely affected both science and public health. Acts of plagiarism have destroyed many scientific careers. Consequently, we want to make clear to you what plagiarism is and penalize acts of plagiarism in a manner that makes clear its seriousness.

Your laboratory reports contain information about the purpose, theory and results of your experiments. Each of you prepares a laboratory report associated only with your name. By implication you are the sole author of that report, and no section of your report can be identical (or nearly identical) to that of another person without attribution. Reports or

sections of reports identical to any other source whether that source be another student, a section of a book, or information obtained from others on the web is treated as plagiarism. In CHM 335 the first instance of plagiarized sections are to receive a grade of 0. For repeat instances of plagiarism, the entire report will receive a 0. Any clear incidence of plagiarism will be reported to the Dean of the appropriate college, the Dean of Arts and Sciences and the Chemistry Department Chair.

In CHM 335 there is one exception to this policy. You often generate your data with a laboratory partner. The original data included in your reports, and only the original data, should be identical to that of your laboratory partners. The other sections of your reports, including all written work and all calculations cannot be identical to anyone including your laboratory partners.

As an example in the second experiment of the semester you determine the diffusive behavior of ionic salts in water. The data from this experiment consists of a sheet of graph paper with multiple curves in different colors generated as a function of time. At the end of the experiment each set of partners has one sheet of graph paper with the same curves. The data is then photocopied in color, and each partner receives either the original sheet or the photocopy. This original data should be included in each partner's laboratory report, and the data are clearly identical for each partner.

To analyze this data each partner must separately generate numbers from the graphs using a ruler and a set of points on the graph chosen by each individual in a manner to be discussed later this semester. It is impossible for the generated numbers to be identical between partners even though the curves are identical. Identical generated data from this experiment or any other experiment is an example of plagiarism.

Another possible source of plagiarism can occur if you do not analyze your data by yourself. Students often work together, and the plagiarism policy is not designed to discourage collaborative learning. However, in the end, your calculations must be your own done in your unique fashion. The sections of your reports containing the calculations must not be identical or nearly identical to anyone else. From experience it is not possible for any two people analyzing the same data to obtain exactly the same set of calculations in the same order with the same final results. To avoid even the appearance of plagiarism, if you decide to study with another student, you must perform your calculations by yourself or with the help of one of the instructors. Nearly identical calculation sections are examples of plagiarism.

11 Schedule

1. Week 1: Lecture and Orientation (Meet in 472 Beaupre)
Tuesday - Sept. 12
Wednesday - Sept. 6
Thursday - Sept. 7
Friday - Sept. 8
2. Week 2: Spectrum of a Conjugated Dye

- Tuesday - Sept. 19
 Wednesday - Sept. 13
 Thursday - Sept. 14
 Friday - Sept. 15
3. Week 3: The Diffusion of Salt Solutions
 Tuesday - Sept. 26
 Wednesday - Sept. 20
 Thursday - Sept. 21
 Friday - Sept. 22
4. Week 4: The Viscosity of Polyvinyl Alcohol, Part a (Meet in 472 Beaupre for a preliminary lecture)
 Tuesday - Oct. 3
 Wednesday - Sept. 27
 Thursday - Sept. 28
 Friday - Sept. 29
5. Week 5: The Viscosity of Polyvinyl Alcohol, Part b
 Tuesday - Oct. 10
 Wednesday - Oct. 4
 Thursday - Oct. 5
 Friday - Oct. 6
6. Week 6: Makeup
 Tuesday - Oct. 17
 Wednesday - Oct. 11
 Thursday - Oct. 12
 Friday - Oct. 13
7. Week 7: Group A - Heat Capacity Ratio of Gases
 Group B - Vapor Pressure
 Group C - Surface Tension
 Group D - Phase Transitions in Ammonium Chloride
 Tuesday - Oct. 24
 Wednesday - Oct. 18
 Thursday - Oct. 19
 Friday - Oct. 20
8. Week 8: Group B - Heat Capacity Ratio of Gases
 Group C - Vapor Pressure
 Group D - Surface Tension
 Group A - Phase Transitions in Ammonium Chloride
 Tuesday - Oct. 31

- Wednesday - Oct. 25
Thursday - Oct. 26
Friday - Oct. 27
9. Week 9: Group C - Heat Capacity Ratio of Gases
Group D - Vapor Pressure
Group A - Surface Tension
Group B - Phase Transitions in Ammonium Chloride
Tuesday - Nov. 7
Wednesday - Nov. 1
Thursday - Nov. 2
Friday - Nov. 3
10. Week 10: Group D - Heat Capacity Ratio of Gases
Group A - Vapor Pressure
Group B - Surface Tension
Group C - Phase Transitions in Ammonium Chloride
Tuesday - Nov. 14
Wednesday - Nov. 8
Thursday - Nov. 9
Friday - Nov. 10
11. Week 11: Experiment 9, Partial Molar Volume
Tuesday - Nov. 21
Wednesday - Nov. 22
Thursday - Nov. 16
Friday - Nov. 17
12. Week 12: Make up
Tuesday - Nov. 28
Wednesday - Nov. 29
Thursday - Nov. 30
Friday - Dec. 1
13. Week 13: In Class Report Writing Session for Experiment 9 (472 Beaupre)
Tuesday - Dec. 5
Wednesday - Dec. 6
Thursday - Dec. 7
Friday - Dec. 8

12 The CHM 335 Web page

It is expected that for most of you, success in this course will require some level of help beyond classroom instruction. Because some of you may find it difficult to come to the scheduled

office hours, we have installed web pages, including a page that can be used to submit questions. Our course web page can be found at <http://www.chm.uri.edu/courses/?chm335&1>. Questions are submitted by anyone in the class by filling out a form on the web page, and answers are distributed either to the entire class or only to the person asking the question. If the entire class is to receive a copy of the question and answer, the question is treated as anonymous; i.e. the person who asks the question is never identified. In fact, it is possible to submit a question so that even the instructor does not know who submitted the question. Anonymous questions and responses by the instructor are distributed automatically to everyone who has submitted their e-mail address to the instructor. With ordinary electronic mail, there is a private correspondence between the student and instructor. By using the web page, the entire class has an opportunity to learn from the questions submitted.

The use of the web page does not preclude personal interaction between any of you and the course instructor. Dr. Freeman, Dr. Yang and the teaching assistants all have regular office hours, and you are all encouraged to make use of these hours. Alternate meeting times can be arranged by appointment. Additionally, you can contact the instructors by e-mail or telephone. The e-mail addresses and phone numbers for all the instructors are given in this syllabus.

To receive copies of the submitted questions and the answers to the questions, you must submit your e-mail address. There are two methods for subscribing to the CHM 335 list. In the first method send your e-mail address to freeman@chm.uri.edu. The subject of the message should be "CHM 335 e-mail list," and the text should contain the e-mail address you would like to use. It is important to identify yourself as a student in CHM 335, so that your name is not added to the list associated with another course. The second method uses our web page. Go to our home page (<http://www.chm.uri.edu/courses/?chm335&1>) and click on "Subscribe to the CHM 335 list." On the resulting form, enter your e-mail address, click on the small "subscribe" button and then click on the submit button. You can also use this form to unsubscribe from the list in case you drop CHM 335.

Any student in CHM 335 can submit questions and comments to the instructors. Submission of such comments or questions must be made using the WWW home page for this course. The address (URL) of our home page is <http://www.chm.uri.edu/courses/?chm335&1>. To submit a question to the list, you must click on the highlighted text that says "submit a question to the CHM 335 list." As an example of how to use the list, suppose a student in our class, Ms. Benzene Ring, wonders, "How do I propagate errors for $\ln p$ where p is the pressure?" (If you don't know what this means, don't worry. You will understand the question later in the semester). To obtain an answer to her question, Ms. Ring links her web browser (e.g. Netscape or Microsoft Internet Explorer) to <http://www.chm.uri.edu/courses/?chm335&1>, and she then clicks on the text linking her to the page for questions (i.e. the highlighted text that says "submit a question to the CHM 335 list"). Ms. Ring then enters her e-mail address in the appropriate box and specifies whether she wants her question to be answered to the entire CHM 335 class or to her alone. Ms. Ring then types in the large box

How do I propagate errors for $\ln p$ where p is the pressure?

Ms. Ring then clicks the “send” button. Ms. Ring’s question is received by the course instructors. One of the instructors then sends an e-mail message to the whole list that might be

Subject: errors

The question is: How do I propagate errors for $\ln p$ where p is the pressure?

Answer: Take the differential. Let $y = \ln p$. Then $dy = dp/p$, and the error in y ($\ln p$) is given by $e(y) = e(p)/p$ where $e(y)$ means the error in y .

Now Ms. Ring and the entire class have an answer to her question.

If the answer to the question can be sent to the entire list, the answer will not indicate who asked the question. If Ms Ring wants to ask the question with full anonymity so that even the instructors have no idea who asked the question, the e-mail portion of the form can be left blank. Of course, if the e-mail section of the form is blank, the answer must be sent to the list and not just to the sender.

Because many questions may contain mathematical formulas, we need a notation to communicate the special symbols used in the course. To avoid confusion, it is most useful if we agree on the same set of symbols. The symbols that follow are taken from a language called \LaTeX . \LaTeX is a language that is frequently used to prepare scientific documents, and \LaTeX can be used to translate special symbols into simple text characters. By learning \LaTeX notation, you will learn a widely used method to communicate mathematical symbols via e-mail. The instructor plans to use these symbols in answering your questions, and it is asked that you use the same symbols in posing questions. The most important symbols are the following:

1. Greek letters are represented by \backslash followed by the name of the letter. For example α is typed $\backslash\alpha$, β is typed $\backslash\beta$, and so on. A Greek letter is made upper case by making the first letter of its name upper case. For example, the letter Δ is typed $\backslash\Delta$.
2. Subscripts are represented by $_{} \{ \}$ where the brackets contain the subscripts. For example, μ_{ij} is typed $\backslash\mu_{} \{ ij \}$.
3. Superscripts are represented by $\hat{} \{ \}$ where the brackets contain the superscripts. For example, β^{12} is typed $\backslash\beta\hat{} \{ 12 \}$.
4. Infinity (∞), is typed $\backslash\infty$.
5. The integral sign \int is typed $\backslash\int$. The limits on a definite integral are included by introducing subscripts and superscripts. As an example $\int_0^\infty e^{-x^2} dx$ is typed $\backslash\int_{} \{ 0 \} \hat{} \{ \infty \} e\hat{} \{ -x\{ 2 \} \} dx$.
6. The partial derivative symbol ∂ is typed $\backslash\partial$.

7. The summation sign \sum is typed `\sum`. The lower and upper limits of summation are included as subscripts and superscripts. As an example $\sum_{n=0}^{\infty} 1/n^2$ is typed `\sum_{n=0}{\infty} 1/n{2}`.
8. Square roots $\sqrt{a+b}$ are typed `\sqrt{a+b}`.
9. The arrow in chemical reactions \rightarrow is typed `---`. For example $C+O_2 \rightarrow CO_2$ is typed `C + O_{2} ---> CO_{2}`.

Let us now look at another example of a question submitted using the web. In this case, Ms. Ring has a question requiring an equation. This might be a real question. If you don't understand the context, don't worry. You will understand the details of the question later in the course. Suppose Ms. Ring wants to ask

"In deriving the expression for the phase equilibrium line between gas and liquid, when evaluating the integral expression

$$\ln \frac{p_2}{p_1} = \int_{T_1}^{T_2} \frac{\Delta H}{R} \frac{dT}{T^2}$$

ΔH is taken outside the integral. What is the justification for this?"

To submit the question, Ms. Ring uses her web browser to attach to <http://www.chm.uri.edu/courses/?chm335&1>, clicks on the line that says, "submit a question to the CHM 335 list," and then Ms. Ring enters the information requested by the form. If Ms. Ring wishes to remain anonymous, Ms. Ring leaves the e-mail box blank. Ms. Ring then types into the large box

In deriving the expression for the phase equilibrium line between gas and liquid, when evaluating the integral expression

$$\ln p_2/p_1 = \int_{T_1}^{T_2} (\Delta H/R) dT/T^2$$

ΔH is taken outside the integral. What is the justification for this? and clicks on the submit button. Ms. Ring's question is received by the course instructors. The answer will be sent either to Ms. Ring alone, or preferably to the entire class if the appropriate box is checked. One of the instructors might reply

Subject: Experiment Number 13 Question

The question is: In deriving the expression for the phase equilibrium line between gas and liquid, when evaluating the integral expression

$$\ln p_2/p_1 = \int_{T_1}^{T_2} (\Delta H/R) dT/T^2$$

ΔH is taken outside the integral. What is the justification for this?

The answer is: ΔH is only weakly dependent on temperature. Then enthalpy change can be taken outside the integral to a good approximation.

Remember, your first task is to subscribe to the CHM 335 list either by e-mail or filling out the form on our web pages. You can then send questions and comments to the course

instructors using your web browser starting at the URL
<http://www.chm.uri.edu/courses/?chm335&1>

13 First Two Laboratory Projects

13.1 Experiment 10(34) - Absorption Spectrum of a Conjugated Dye

In general chemistry you have learned that the electrons in atoms and molecules occupy discrete energy states called energy levels. The electrons fill the energy levels sequentially (called an *aufbau*) according to the Pauli Exclusion Principle where no more than two electrons can fill any one energy level. If two electrons occupy the same energy level, those two electrons must have the opposite spin. With the exception of the hydrogen atom, the values for the energy levels available to the electrons have no simple, closed-form expression. Consequently, it is useful to use model systems to describe the energy states.

A good example of molecules whose energy levels can be represented by a simple model are linear polyenes. The π -electrons in such systems can be modeled by the energy levels of electrons confined to move freely in a one-dimensional box. The energy levels for the particle in a box can be found analytically as explained in CHM 432, and the resultant expression is

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}, \quad (1)$$

where L is the box length, m is the particle mass (the mass of an electron in our problem), $\hbar = h/2\pi$ with h Planck's constant (6.626×10^{-34} Js), and n is called the quantum number and can take on values = 1,2,3,... The electrons in linear polyenes can be imagined to occupy electronic energy levels with values given by Eq.(1). When using this particle in a box model, the box length is taken to be the linear distance the delocalized π -electrons can move in the molecule.

Electronic energy states in atoms and molecules are often probed experimentally using spectroscopy. In a spectroscopic experiment photons (particles of light) are either absorbed or emitted as the electrons "jump" between the available energy levels. In this experiment where absorption spectroscopy is used, the energy of the absorbed photons must equal the difference in energy between the electron in the highest occupied energy level and the lowest unoccupied energy level. The purpose of this experiment is to measure the absorption spectrum for a series of linear polyenes and determine the extent to which the particle in a box model describes the experimental data.

13.1.1 Reading Assignment

Textbook, Experiment number 10(34), pages 117-122. The actual procedure that is followed in this experiment differs from the procedure outlined in the textbook. Your TA will provide an explanation of the modifications.

13.1.2 Experimental Notes

1. Be sure to prepare the solutions in the fume hood and use rubber gloves. The dyes used in this experiment are toxic.
2. Your TA will give you detailed instructions about how to use the spectrometer.

13.1.3 Laboratory Reports

Your reports should include the following sections:

1. Cover Page: Include the title of the experiment, the date the experiment was performed and the name of your laboratory partner.
2. Introduction: You should discuss the purpose of the experiment and give an overview of the key ideas. Make clear how the measurement of the absorption maximum enables the determination of information about the structure of the molecules.
3. Theory: Discuss the particle in a box model, and use the model to derive the key formulas used to analyze the data.
4. Procedure: Attach the outline of the procedure that your TA's signed in class.
5. Original Data: Attach the original data signed and dated by your TA.
6. Calculations:
 - (a) Use Eq. (6) in the textbook to calculate λ_{max} for each dye studied. Choose α to match the experimental results for the dye specified by your TA. Then using this value of α , determine λ_{max} for the remaining dyes.
 - (b) Prepare a table that lists each compound, the value of p for the compound, and a comparison of the calculated values of λ_{max} with the experimental ones.
7. Summary and Discussion: Briefly summarize the experiment and your findings. Include all key numerical results including the value of α that you determined. Discuss the extent to which the particle in a box model is appropriate for the set of molecules studied in this experiment.

13.2 The Diffusion of Salt Solutions into Pure Water

When a drop of ink is placed in a beaker of water, the ink is observed to migrate from the initial location of the ink drop to the remainder of the water. This migration is called *diffusion* and continues until the ink concentration is uniform throughout the solution. In this experiment we study the time dependence of one property of the diffusion process.

Instead of a drop of ink, you use an ionic solution and let it diffuse into pure water to determine how the ionic concentration gradient depends on time. You begin with a

rectangular cell with the bottom half filled with the salt solution and the top half filled with water as depicted in Figure 1. As time passes, the salt in the concentrated layer migrates to regions of low salt concentration. If the salt concentration c is measured as a function of time t and coordinate x [See Figure 1 for the definition of x], the concentration profile appears as in Figure 2 [c_0 is the initial salt concentration in the salt rich layer]. The two lines drawn in Figure 2 represent concentration profiles at two different times. Can you tell which of the two times is earlier? Using Fick's Law of Diffusion [See Engel and Reid, Sections 17.2 through 17.4], it is possible to show that

$$\frac{\partial c}{\partial x} = Ae^{-x^2/2\sigma^2(t)} \quad (2)$$

where A is a constant and σ is called the *standard deviation*. Equation (2) is often called a Gaussian distribution. A graph of Eq.(2) is shown in Figure 3. The standard deviation σ is a measure of the width of the distribution. It is possible to show that the area under a Gaussian between $-\sigma \leq x \leq \sigma$ comprises $\sim 68\%$ of the total area under the Gaussian curve. Therefore, if σ is large [the light line in Figure 3], the distribution is wide and if σ is small [the dark line in Figure 3] the distribution is narrow.

In this experiment you measure σ directly in a diffusion cell as a function of time. You can be expected to find that

$$\sigma = Kt^\alpha \quad (3)$$

where K and α are constants. Relations of the form of Eq. (3) are often called *power laws*. You should find that α is particularly simple.

The method you will use to determine the concentration gradient of the solution as a function of time is based on the refraction of light. When light travels from one medium to another, the light bends. This bending of light is called *refraction*. In common experience, if a spoon is placed in a glass of water, it appears bent because of refraction. The amount that light is bent is proportional to a property of a medium called the *index of refraction* n . In a salt solution, the index of refraction depends on the type of ionic species and the concentration of the solution. In our cell where there is a concentration gradient, the index of refraction depends on the coordinate x . Since the index of refraction is proportional to the concentration, it is easy to see that

$$\frac{\partial c}{\partial x} = K' \frac{\partial n}{\partial x} \quad (4)$$

where K' is a constant of proportionality. We can then measure the concentration gradient by measuring the gradient of the index of refraction.

In this experiment the index of refraction gradient is measured by passing a laser light beam perpendicular to a glass stirring rod that acts as a cylindrical lens. The stirring rod is set at an angle of about 45 degrees to the horizon. The resulting light image is projected onto the surface of the diffusion cell. The line is at an angle of approximately 45 degrees to the horizontal axis of the cell and perpendicular to the direction of the glass rod. If the line of light is passed through a solution of uniform concentration, the image on the

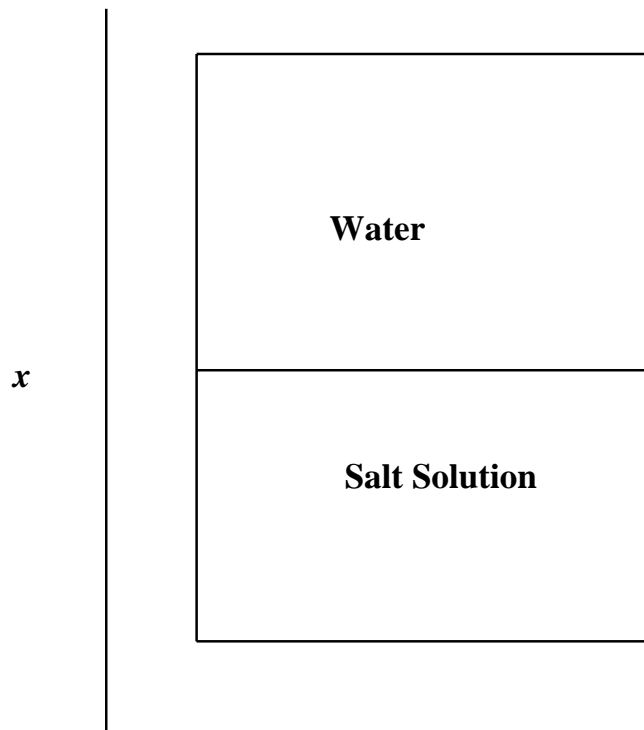


Figure 1

Figure 2

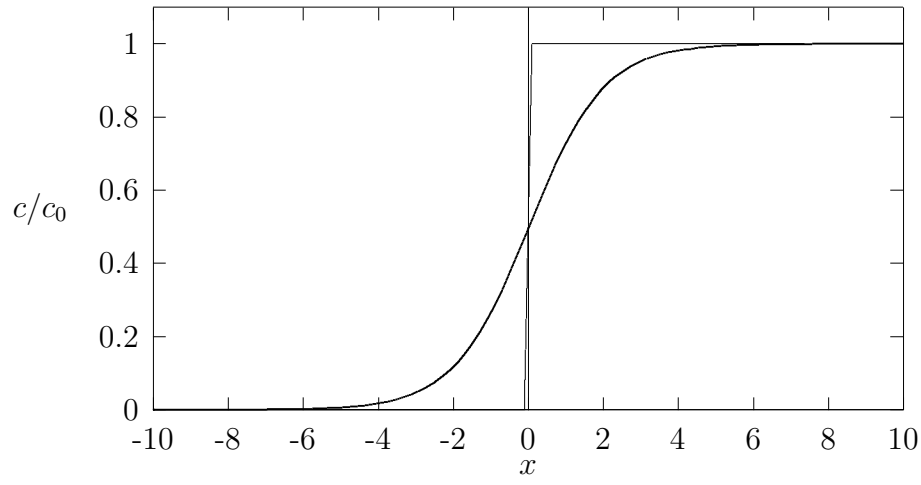
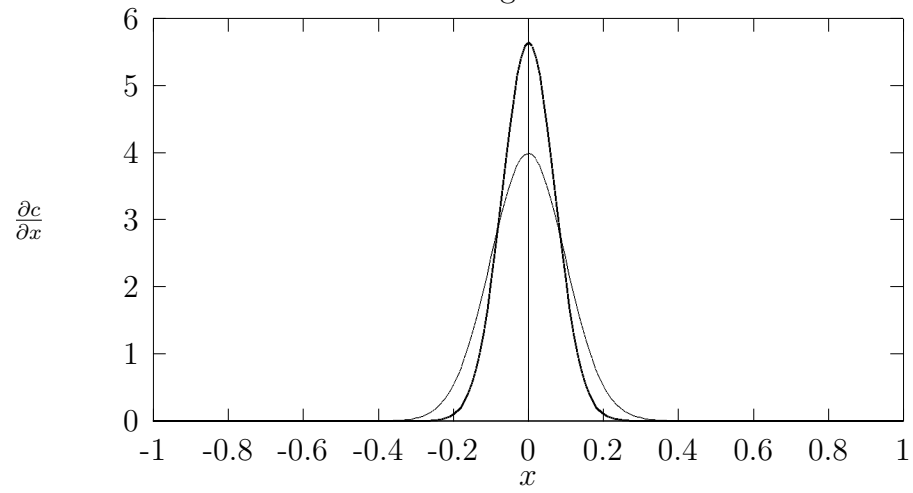


Figure 3

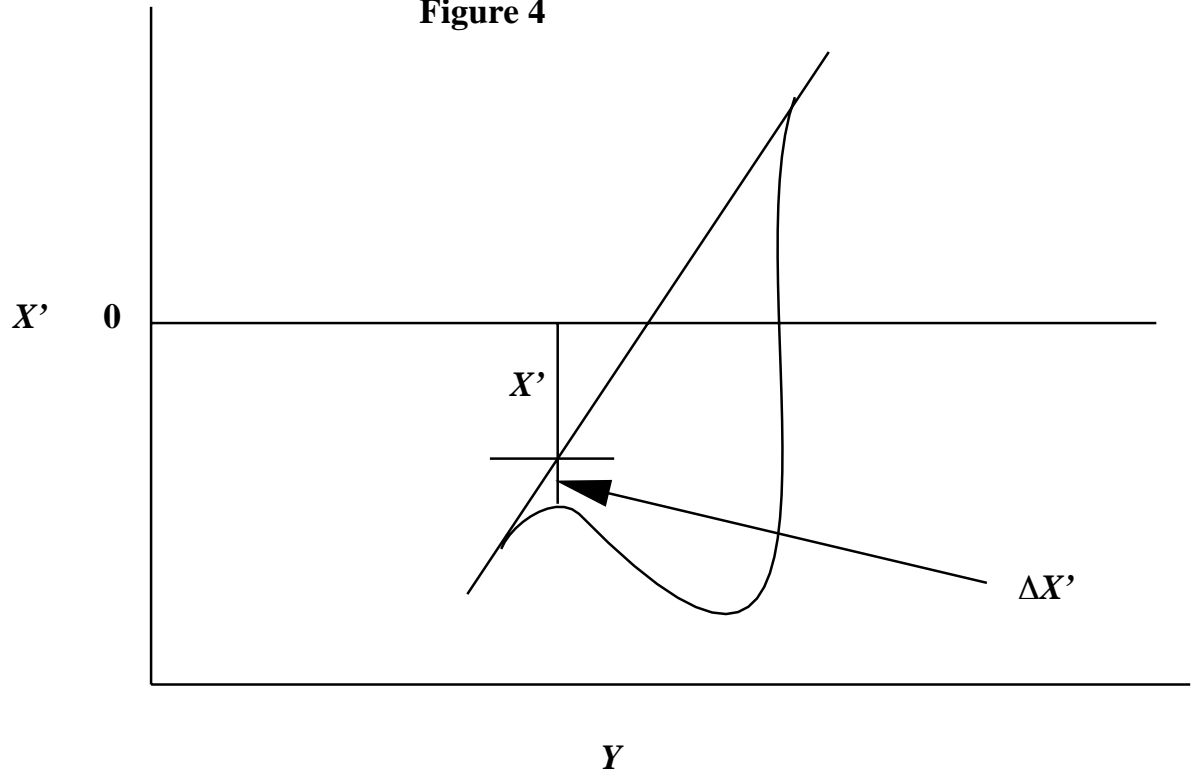


other side of the diffusion cell, that we project on to a screen, is parallel and longer than the incoming image. However, if the solution has a concentration gradient, the component of the incoming light in the x direction deflects by differing amounts depending on the concentration gradient. The resulting projected image appears as in Figure 4. In Figure 4, the straight line represents the image that appears in a solution of uniform concentration, and the curved line represents the image that appears if the solution has a concentration gradient. At any location along the abscissa, X' in the graph is proportional to the coordinate x in the cell and the difference $\Delta X'$ is proportional to the concentration gradient. A plot of $\Delta X'$ as a function of X' then is a Gaussian distribution. The dependence of the standard deviation of the resulting Gaussian on time is the same as the dependence of the standard deviation in Eq. (2) on time. Consequently, we can work with the plotted data directly.

13.2.1 Procedure

1. Make up 2 M solution of the salt given to you. Record in your laboratory notebook the name of the salt. Be sure to stir the solution thoroughly to completely dissolve the salt in water.
2. You begin this experiment by solving a simple riddle qualitatively. Fill a 50 ml beaker with distilled water and another 50 ml beaker with your salt solution. Imagine that you have no idea which beaker contains water and which contains the salt solution. Given only a pipette and pipette bulb, you are to determine which beaker contains the water and which contains the salt solution.
3. For the main part of this laboratory, the experimental components consist of a low power helium-neon laser, a cylindrical lens made from a glass rod, a cuvette (diffusion cell) on a laboratory jack, and a a piece of graph paper that is to be used as a screen. The glass rod, the cuvette and the screen are mounted perpendicular to the laser beam as shown in Figure 5. To scan several depths in the cuvette simultaneously, the beam of light from the laser is fanned through the glass stirring rod mounted at about 45 degrees to the horizontal. The rod is placed so that its edge is a distance of approximately 40 cm. from the center of the cuvette. The cuvette is mounted on a laboratory jack to facilitate vertical positioning. The graph paper screen is placed on the wall at the end of the lab bench.
4. Fill the cell halfway with the salt solution.
5. Raise or lower the cell until the laser hits the meniscus. A vertical line will appear on the screen. Trace this line.
6. Slide the ring stand with the glass rod attached back so that the rod is between the cuvette and the laser making a line at a 45 degree angle. Trace the diagonal laser line on the graph paper.

Figure 4



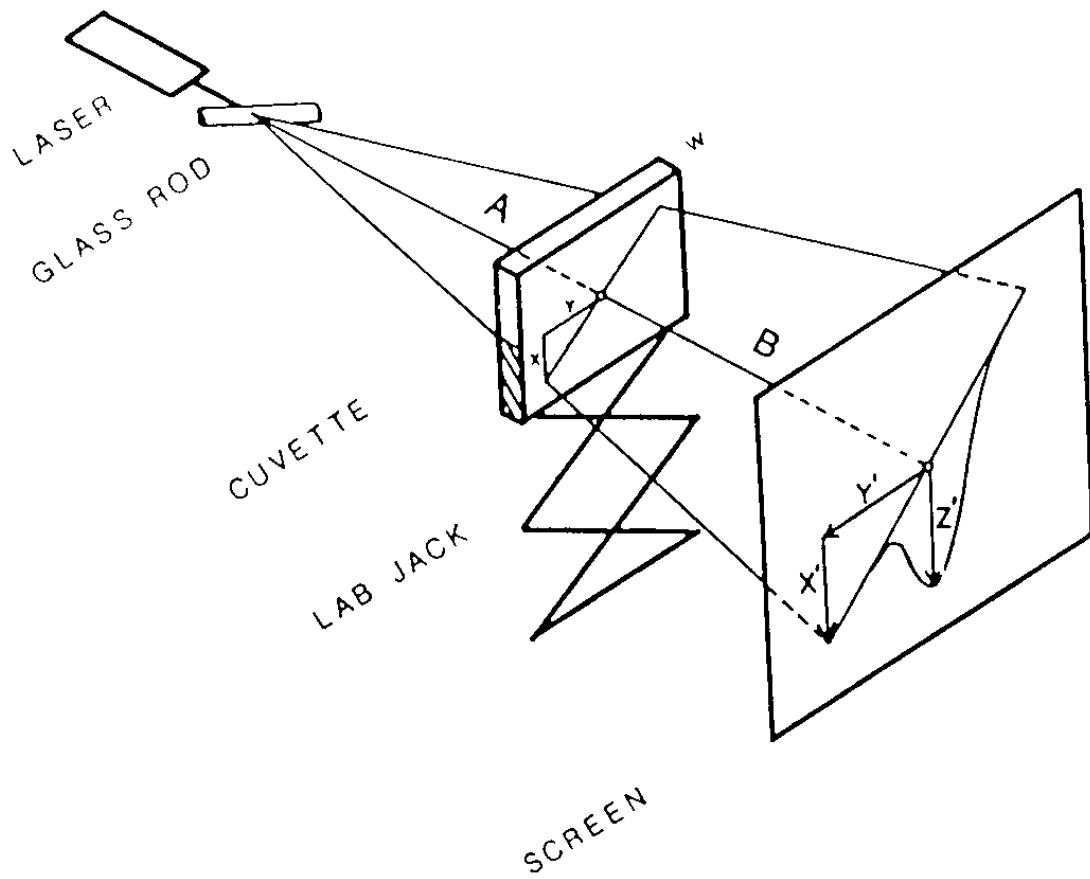


Figure 3. Line drawing of the experimental apparatus.

7. Add water to the cuvette by using a pipette and a floating cork. Make sure the water is added slowly and directly on the cork so as not to disturb the interface. This process should take a few minutes. You should start your stopwatches at the time you begin to add water to the cuvette.
8. When you have waited sufficient time so that a stable curve appears entirely on your graph paper, trace the curve on the graph paper and note the time.
9. Trace new lines at 5 minute intervals for 40 minutes.

13.2.2 Laboratory Report

Your reports should consist of the following parts:

1. A cover page giving the date the experiment was performed, the name of your laboratory partner and the title of the experiment;
2. Your procedure signed by your TA;
3. A description of how you solved the riddle at the beginning of the experiment. Be sure to include a description of the physical principals used to solve the riddle. These principals, of course, should relate to the rest of the experiment.
4. The raw data [in this case the original graph paper screen signed by your TA];
5. A table of 10 values of X' and $\Delta X'$ calculated from the graph paper for each time;
6. Graphs of $\Delta X'$ as a function of X' for each time;
7. A statement of the precise definition you used for the width w of each graph;
8. A table of w , t , $\ln w$ and $\ln t$;
9. Plot the values of $\ln w$ as a function of $\ln t$ using the calculated points;
10. Draw a straight line that is your best estimate of the linear relationship between $\ln w$ and $\ln t$;
11. Determine the slope of your graph;
12. Given that the slope is either an integer or a half integer, round off the value of your slope to the nearest half integer;
13. Express the dependence of the the width as a power law. It is important to recognize that the width w is proportional to the standard deviation σ of the Gaussian curves.
14. Summary and Discussion: Briefly summarize the experiment and your findings. Include all key numerical results; e.g. power law. Examine Eq. (17.18) of the CHM 431 textbook, and compare your measured value of α with the power law of the mean squared displacement in diffusion processes as a function of time.

14 Second Two Laboratory Experiments

14.1 Notes on Error Analysis

14.1.1 Reading Assignment

Text pages 3-6, 17-22, 26-41.

14.1.2 Introduction

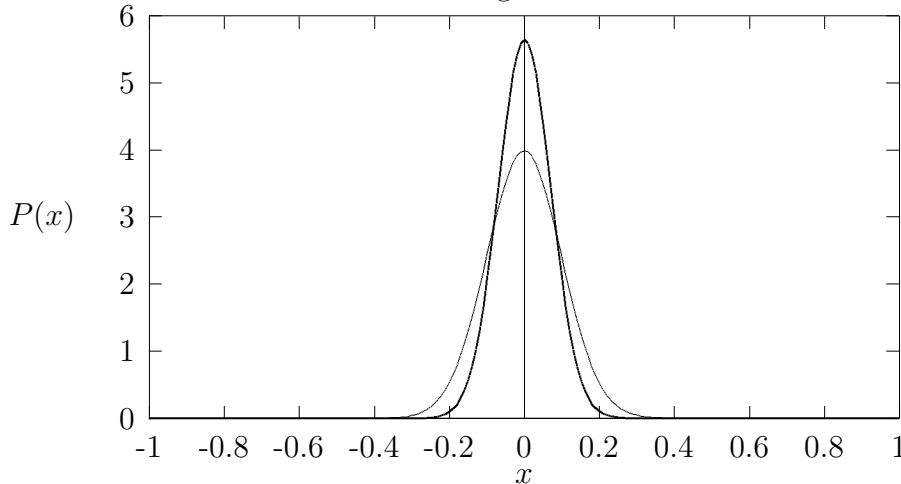
The study of the analysis of experimental error is an important goal of CHM 335. The instructors of this course feel error analysis to be sufficiently important that from 30% to 40% of the grade on your later laboratory reports are based on this topic. We encourage you to learn this topic well as early in the semester as possible. Many aspects of error analysis are complex, and the lecture together with these course notes and the appropriate sections of the text book should be helpful. In preparing your reports we encourage you to come to the instructors for help with error analysis to make sure you understand the key ideas.

To illustrate the importance of the analysis of errors, suppose a surveyor measures the distance between Kingston and New York City, and reports the distance to be 150. miles. While stating the distance to be 150 miles might be sufficient for the Auto Club, from the point of view of the physical sciences the reported distance alone is inadequate. For example, the surveyor might mean 150 miles \pm 1 mile or the surveyor might mean 150 miles \pm 100 miles. The first possibility tells us that it should take about 3 hours to drive to New York while the latter is only a rough order of magnitude estimate. While the auto club probably assumes something like 150 miles \pm 1 mile, anyone in the sciences must provide not only the result of a measurement, but the extent to which the result is known as well. Starting with the third experiment when you report an experimental result, you must include both the value obtained as well as the error that provides information about how well you know the result reported.

14.1.3 Systematic and Random Errors

Experimental errors are often divided into two categories. The first category, usually called *systematic errors*, refers to mistakes inherent in a particular apparatus. For example, suppose we measure the length of the cover of your textbook with a metal ruler. The total length of the ruler expands and contracts depending on the temperature of the room. The ruler might be calibrated assuming the temperature is 20°C, and if the actual room temperature is 25°C, the total length of the ruler must be longer than originally measured. Consequently, any measurement of length at the higher temperature can be expected to give a result that is too short. It is important to recognize that systematic errors have a definite algebraic sign. In the case of our expanded ruler, the error is necessarily negative. There is some control over such systematic errors. For example, if the coefficient of thermal expansion of the ruler is known, the measured length can be corrected to provide the true length. Associated with

Figure 6



systematic errors is the term *accuracy*. An apparatus is said to be accurate to the extent that the systematic errors are minimized.

The second category of error is associated with inherent random fluctuations of any measurement apparatus, and consequently we refer to such errors as *random errors*. Again, we can take the measurement of the length of the cover of our textbook as an example. If we measure the length of the book many times, we can expect to obtain a range of values that fluctuates about a mean value. The fluctuations come from many sources. For example, you might not place the bottom of the ruler at exactly the same place for each measurement. The thermal fluctuations of the atoms in the ruler change its length slightly providing another source. The temperature of the room also fluctuates changing the length of the ruler in a random way. Such random fluctuations are inherent, and cannot be eliminated. By hypothesis, we assume the correct result (in the sense of random errors) for some measurement can be obtained from the arithmetic mean of an infinite set of measurements. We control random errors by increasing the number of measurements. Unlike systematic errors, random errors have no definite algebraic sign, and we denote the size of random errors with \pm notation. We use the term *precision* to express the magnitude of the random errors. We call a measurement precise if its random errors are small.

14.1.4 The Distribution of Random Errors - The Gaussian Distribution

An important theorem in probability theory, called *the central limit theorem*, states that random, independent measurements of a quantity (call it x) must be distributed according to a *Gaussian* distribution that takes the form

$$P(x) = \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left\{-\frac{(x - \mu)^2}{2\sigma^2}\right\}. \quad (5)$$

The meaning of the parameters σ and μ are given shortly. Equation (5) is plotted in Figure 6 for the case that $\mu = 0$ and $\sigma = 1$. (the broader peak) and $\mu = 0, \sigma = 0.5$ (the narrower peak). The probability function $P(x)$ is normalized in the sense that

$$\int_{-\infty}^{\infty} dx P(x) = 1 \quad (6)$$

implying that $P(x)$ is a well-defined probability function. The average of x is given by

$$\langle x \rangle = \int_{-\infty}^{\infty} dx x P(x) = \mu \quad (7)$$

so that the point $x = \mu$ locates the peak of $P(x)$ and, by symmetry, is also the mean. A second important integral is

$$\langle x^2 \rangle - \langle x \rangle^2 = \int_{-\infty}^{\infty} dx x^2 P(x) - \mu^2 = \sigma^2. \quad (8)$$

The parameter σ is often called the *standard deviation* and is a measure of the width of the distribution. It is possible to show that

$$\int_{-\sigma}^{\sigma} dx P(x) \cong 0.66 \quad (9)$$

and

$$\int_{-2\sigma}^{2\sigma} dx P(x) \cong 0.95. \quad (10)$$

Because $P(x)$ represents the distribution of measured values about the mean μ , Eqs. (9) and (10) imply that the result of any measurement has a 66% chance of occurring in the range $\mu \pm \sigma$ and a 95% chance of occurring in the range $\mu \pm 2\sigma$.

14.1.5 Calculation of Random Errors for a Finite Set of Laboratory Measurements

The distribution $P(x)$ is exact only for an infinite set of independent measurements. Naturally, in a laboratory we can only generate a finite set of measured data. For a finite set of measurements presumably distributed according to $P(x)$, we estimate μ and σ . To estimate σ and μ for a finite set of measurements, we imagine the set of N measured values to be x_1, x_2, \dots, x_N . The mean is given by the standard expression

$$\mu = \frac{x_1 + x_2 + \dots + x_N}{N} = \frac{1}{N} \sum_{i=1}^N x_i. \quad (11)$$

To determine the standard deviation we use

$$\langle x^2 \rangle = \frac{1}{N} \sum_{i=1}^N x_i^2, \quad (12)$$

so that

$$\sigma^2 = \frac{1}{N} \sum_{i=1}^N x_i^2 - \left(\frac{1}{N} \sum_{i=1}^N x_i \right)^2. \quad (13)$$

The standard deviation σ represents a measure of the fluctuations of x from a large set of measurements. Although σ is of interest, our principle concern is the fluctuations of the calculated mean. It is not hard to show that the standard deviation of the mean itself σ_N is given by

$$\sigma_N^2 = \frac{\sigma^2}{N-1} \quad (14)$$

(the $N-1$ rather than N comes from the loss of one degree of freedom in the estimation of the mean). We then report our results in the form

$$x = \mu \pm 2\sigma_N. \quad (15)$$

From the discussion at the end of Section 14.1.4, there is a 95% probability that the correct result is in the range given in Eq.(15).

Exercise: A student measures the mass of an iron nail with a balance and obtains the following results: 5.500 g, 5.560g, 5.550g, 5.550g and 5.590g. Calculate the result that the student should report.

Answer: The average mass is found to be 5.550 grams, $\sigma = .02898$ grams, and $\sigma_N = .01449$ grams. The reported result should be $5.55 \pm .03$ grams. Notice the first significant digit in the error defines the number of significant figures that should be reported for the average. We don't use simple first year chemistry rules for significant figures, but rather we use the computed error to define the number of significant figures to be reported.

The answer to the previous exercise requires rounding off to the last significant digit. The usual rule for rounding off digits is to increase by one digit if the next digit is greater than 5, do not increase by one digit if the next digit is less than 5, and increase by one digit if the next digit is 5 and the result of the round off is even. For example 5.56 becomes 5.6, 5.54 becomes 5.5 and 5.55 becomes 5.6. However 5.65 becomes 5.6.

In most cases it is important to repeat a measurement about 10 times to ensure the generation of proper statistics so that the magnitude of the random errors can be estimated. With the exception of the third experiment, in CHM 335 there is not adequate time to perform 10 measurements. Consequently, the size of the random errors for a particular measurement is estimated. As a rule of thumb the size of the random error for a measurement is ± 2 in the last reported digit(we assume the ± 2 error to be at the 95% confidence level; i.e. 2 standard deviations of the mean). For example, when using the analytical balances, it is possible to read the mass to four decimal places. If a particular sample has a measured mass of 4.2345 grams, it is safe to assume the mass to be $4.2345 \pm .0002$ grams. In the case of volumetric flasks, the size of the random error is printed on the flask below the label of

the total volume of the flask. When using a volumetric flask, it is important to record the error in the flask volume when recording your data. For volumetric pipettes (so called class A glassware), the random errors are those given in the following table:

Volume (ml)	random error (ml)
1.	0.01
2.	0.01
5.	0.02
10.	0.02
20.	0.03
25.	0.03
50.	0.04
100.	0.06

14.1.6 Error Propagation

1. Treatment of Systematic Errors

Perhaps the most difficult topic in error analysis is how errors in basic measurements propagate to a final result. For example the pH of a solution is defined by

$$pH = -\log_{10}[\text{H}^+] = -\ln[\text{H}^+]/2.3 \quad (16)$$

(really, the definition of the pH is in terms of activities of the hydrogen ion, but we can take it to be defined in terms of concentrations for our purposes). If the hydrogen ion concentration is measured directly with a known statistical error, we now discuss how to determine the error in the pH . The generic form of this problem can be discussed in terms of a function $f(x, y, z, \dots)$ where x, y, z, \dots represent fundamental quantities measured directly, and f represents the actual physical quantity to be reported. In our pH example, there is a single fundamental measured quantity $x = [\text{H}^+]$, and the function is $f(x) = -(1/2.3) \ln x$.

The total differential of a function provides the mathematics necessary to solve the problem of error propagation. The total differential plays a central role in the study of thermodynamics as well, and total differentials are discussed in detail in CHM 431. As a reminder, the total differential of a function f is defined by

$$df = \left(\frac{\partial f}{\partial x}\right)_{y,z,\dots} dx + \left(\frac{\partial f}{\partial y}\right)_{x,z,\dots} dy + \left(\frac{\partial f}{\partial z}\right)_{x,y,\dots} dz + \dots \quad (17)$$

The total differential of f is the infinitesimal change in f that results from the infinitesimal changes in its variables expressed as dx, dy, dz, \dots . In the case of errors, df is the error in f that arises from small errors dx, dy, dz, \dots in each of its variables. The interpretation of df just provided in terms of errors is valid for systematic errors.

We shall discuss the modifications for random errors in Section 2. We first try the application of the total differential to our pH example. Using $f = pH$ and $x = [H^+]$ we have

$$f = - \left(\frac{1}{2.3} \right) \ln x \quad (18)$$

and

$$df = - \frac{dx}{2.3 x}. \quad (19)$$

Then

$$\epsilon(pH) = - \frac{\epsilon([H^+])}{2.3[H^+]} \quad (20)$$

where the notation $\epsilon(x)$ denotes the systematic error in the variable x . The error in the pH is not simply equal to the error in the hydrogen ion concentration itself, but rather the error in the pH is a function of the error in the hydrogen ion concentration as well as other variables. The functional form of the expression for the error cannot be guessed in any simple way, but must be derived from the expression for the total differential. Because systematic errors have a definite algebraic sign, the systematic error in the pH carries sign information.

Let us try another example. The density ρ of a substance is defined by

$$\rho = \frac{m}{V} \quad (21)$$

where m is the mass and V is the volume. For the purpose of this exercise we take the mass and the volume as the quantities actually measured in a laboratory, and we ask how the errors in the mass and volume propagate to give the final error in the density. Writing $\rho = \rho(m, V)$ The total differential of the density is given by

$$d\rho = \left(\frac{\partial \rho}{\partial m} \right)_V dm + \left(\frac{\partial \rho}{\partial V} \right)_m dV \quad (22)$$

$$= \frac{dm}{V} - \frac{m dV}{V^2}. \quad (23)$$

The expression for the total differential in the density can be simplified by dividing the left hand side of Eq.(23) by ρ and the right hand side of Eq.(23) by m/V resulting in the expression

$$\frac{d\rho}{\rho} = \frac{dm}{m} - \frac{dV}{V}. \quad (24)$$

The expression for the propagation of systematic errors then takes the form

$$\frac{\epsilon(\rho)}{\rho} = \frac{\epsilon(m)}{m} - \frac{\epsilon(V)}{V}. \quad (25)$$

The right hand side of Eq.(25) contains sign information. A positive systematic error in the mass results in a positive error in the density whereas a positive systematic error

in the volume results in a negative systematic error in the density. These signs make sense, because increasing the volume decreases the density (the volume term appears in the denominator).

In this course we generally assume that there are no systematic errors. We have included the analysis of this section to provide the needed mathematics for the treatment of random errors. In your laboratory reports you should focus on random errors, the method for which is discussed in the next section.

2. Treatment of Random Errors

In the previous section the expressions for the propagated systematic errors contained sign information. Unlike systematic errors, random errors contain no signs. It is possible to show that the expressions for systematic errors become valid for random errors if each term in the expression derived from the total differential is squared and averaged so that cross terms are eliminated. The proof of this assertion is beyond the scope of this course. To make the procedure explicit, consider a function f of two measured quantities x and y . The total differential of f is given by

$$df = \left(\frac{\partial f}{\partial x} \right)_y dx + \left(\frac{\partial f}{\partial y} \right)_x dy. \quad (26)$$

As before, we replace the differentials by errors

$$\epsilon(f) = \left(\frac{\partial f}{\partial x} \right)_y \epsilon(x) + \left(\frac{\partial f}{\partial y} \right)_x \epsilon(y). \quad (27)$$

Equation(27) is just the expression valid for the propagation of a systematic error. We now square both sides of Eq.(27) and average. By averaging, the cross terms containing terms of the form $\epsilon(x)\epsilon(y)$ average to zero because each individual error can have both positive and negative signs. The final result for the propagated random error is then

$$\epsilon(f)^2 = \left(\frac{\partial f}{\partial x} \right)_y^2 \epsilon(x)^2 + \left(\frac{\partial f}{\partial y} \right)_x^2 \epsilon(y)^2. \quad (28)$$

(We have replaced the symbol ϵ that represents a systematic error by the symbol ϵ to represent a random error.) To apply the procedure to our density example, we can begin with Eq.(25), square both sides and average to obtain

$$\left(\frac{\epsilon(\rho)}{\rho} \right)^2 = \left(\frac{\epsilon(m)}{m} \right)^2 + \left(\frac{\epsilon(V)}{V} \right)^2. \quad (29)$$

Note: Students often mistakenly double a propagated error imagining that the propagated error must be at the single standard deviation (66 %) confidence level. In

fact, the error in the fundamental measured quantities should already be expressed as 2 standard deviations of the mean, and the propagated error represents 2 standard deviations of the mean. No doubling of the propagated error is necessary.

Exercise: An expression for the vapor pressure of a certain liquid from a particular experiment takes the form

$$\ln(p_2 - p_1) = \frac{\lambda}{RT}$$

where p_1 and p_2 are two measured pressures expressed in unitless form, R is the gas constant ($8.3144 \text{ J mol}^{-1}\text{K}^{-1}$), λ is the enthalpy of vaporization, and T is the temperature in degrees Kelvin. Suppose the experiment is used to determine the value of λ , and the measured data are $p_2 = 5.53 \pm 0.02$, $p_1 = 4.07 \pm 0.09$, and $T = 275.1 \pm 0.1\text{K}$. Use the data to determine λ , determine the propagated error in λ and express the result with the appropriate number of significant figures. For simplicity, you should assume R is an exact constant with no associated error (in fact R is an experimental number with associated but small random errors).

Answer:

$$\varepsilon(\lambda) = \sqrt{(R \ln(p_2 - p_1))^2 \varepsilon^2(T) + \left(\frac{RT}{p_2 - p_1}\right)^2 \varepsilon^2(p_1) + \left(\frac{RT}{p_2 - p_1}\right)^2 \varepsilon^2(p_2)}$$

$$\lambda = (9. \pm 1.) \times 10^2 \text{ J mol}^{-1}$$

3. Special Useful Expressions for the Propagated Error

In general to determine the expression for a propagated error, the direct calculation of the total differential is necessary. The previous exercise illustrates the need to actually determine the total differential. However, there are two special cases that are so common that it is useful to derive general expressions for future use. The first is a function that depends *only* on the sums and differences of its variables. For example $f(x, y, z) = x + y - z$ The total differential of f is

$$df = dx + dy - dz. \quad (30)$$

After replacing the differentials by errors and squaring and averaging, the result is

$$\varepsilon(f) = \sqrt{\varepsilon(x)^2 + \varepsilon(y)^2 + \varepsilon(z)^2}. \quad (31)$$

In general the propagated random error for any function of only the sums and differences of its variables is the square root of the sums of the squares of the errors in its variables.

The second important case is a generalization of our density example. Suppose f is a function *only* of products and quotients of its variables. For example take $f = xy/z$. The total differential of f is

$$df = \frac{y}{z}dx + \frac{x}{z}dy - \frac{xy}{z^2}dz. \quad (32)$$

Dividing the left hand side of Eq.(32) by f and the right hand side of Eq.(32) by xy/z , we obtain

$$\frac{df}{f} = \frac{dx}{x} + \frac{dy}{y} - \frac{dz}{z}. \quad (33)$$

We next replace the differentials by errors and square and average to obtain

$$\varepsilon(f) = f \sqrt{\left(\frac{\varepsilon(x)}{x}\right)^2 + \left(\frac{\varepsilon(y)}{y}\right)^2 + \left(\frac{\varepsilon(z)}{z}\right)^2}. \quad (34)$$

We can interpret Eq.(34) by stating the relative error in f is the square root of the sum of the squares of the relative errors in each of its variables. This result can be extended to any number of variables as long as all variables appear as products or quotients only.

14.1.7 Errors from a Graphical Analysis of the Data

Frequently the physical result we seek in an experiment is extracted from the slope or the intercept of a linear graph obtained from a series of data. For example, suppose the goal of an experiment is to determine a value of the gas constant R from pVT data obtained from measurements using one mole of an ideal gas. The value of R can be obtained from the slope of a graph of the product pV plotted as a function of the temperature T . An example of such data is shown in Figure 7. In the data shown in Figure 7, the random errors in the temperature are smaller than the resolution of the graph. The random errors in the pV product are large and depicted by the vertical bar with the horizontal lines at the head and foot associated with each data point. For example the first plotted point in Figure 7 is 22.4 ± 1 l atm. A point is plotted at $pV = 22.1$ l atm and the vertical bar ranges from 21.1 l atm to 23.1 l atm. These vertical bars are called *error bars*. When plotting points containing random errors, associated with each point must be an error bar. There should be error bars parallel to both the x and y -axes, but in this example (and most of the experiments in this course), the error bars for the data plotted on the x -axis are too small to see.

Perhaps the best way to determine the optimal line that connects a set of experimental points is called *linear regression*. Linear regression is discussed in your textbooks, but in CHM 335 we use a simpler method. We draw the best straight line using a visual fit. The

best straight line drawn in this way is depicted as the darkest solid line in Figure 7. From this best fit line, the slope can be determined. The slope so determined is taken to be the slope to be reported in your laboratory experiment. For example, the slope of the best-fit (dark) line in Figure 7 is $0.0814 \text{ l atm K}^{-1}$. We ask that you analyze your data in this way using graph paper rather than a computer. In the final experiment of the year, the reports are written in class, and by preparing your graphs using graph paper rather than a computer, you have the necessary practice for the final laboratory project.

To obtain an estimate of the errors associated with the slope of a graph (or equivalently the intercept), the plotted error bars are used to obtain extreme values of the possible slope and intercept. In this course we take the first and last plotted points, and connect the bottom of the first error bar to the top of the last error bar with a straight line. Similarly, we connect the top of the first error bar with the bottom of the last error bar with another straight line. These two lines are depicted in Figure 7 as the two light solid lines. We call the slope of the light solid line having the largest slope m_+ and the slope of the light solid line having the smallest slope m_- . We estimate the error in the slope to be

$$\varepsilon(m) = \frac{1}{2}(m_+ - m_-) \quad (35)$$

For example, in Figure 7 $m_+ = 0.0942 \text{ l atm K}^{-1}$, $m_- = 0.0784 \text{ l atm K}^{-1}$ so that the error is $\varepsilon(m)=0.008 \text{ l atm K}^{-1}$. We then report the gas constant to be $R = 0.0814 \pm 0.008 \text{ l atm mol}^{-1} \text{ K}^{-1}$. The error estimated in this fashion can often be an over estimate. Consequently we can take the error calculated in Eq.(35) to represent two standard deviations of the mean; i.e. at the 95% confidence level.

Although linear regression is capable of providing a more careful estimate of the error in the slope from a set of data, the method expressed in Eq.(35) is more suitable for the purposes of CHM 335.

The error in the intercept b is treated analogously.

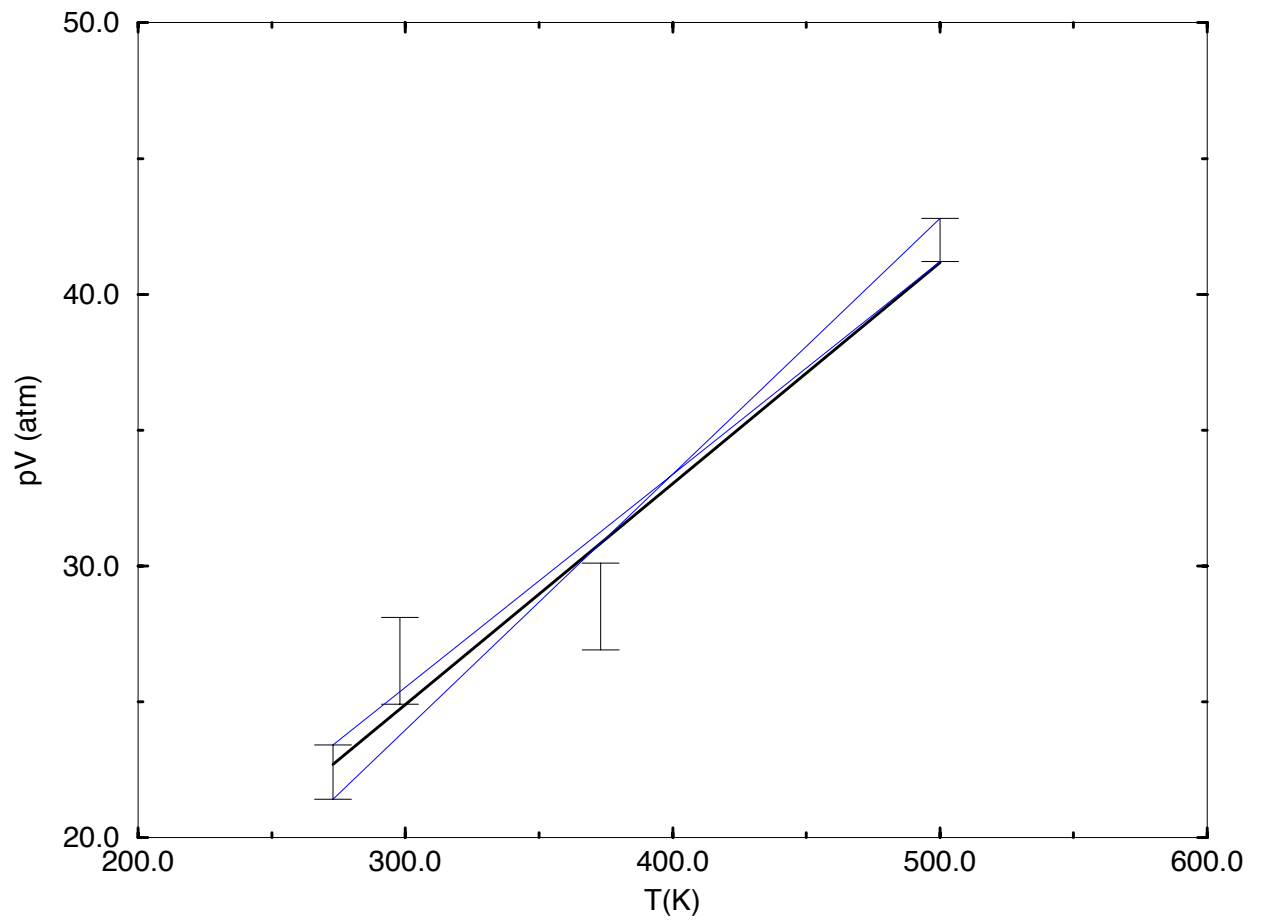
$$\varepsilon(b) = \frac{1}{2}(b_+ - b_-) \quad (36)$$

14.2 Experiment 9(28)a: The Intrinsic Viscosity of Polyvinyl Alcohol

14.2.1 Introduction and Reading Assignments

In this and the next experiment, you use the measurement of the the viscosity of liquids to illustrate two things. In the first experiment you learn the methods of error analysis. In the experiment to be performed in the following week, you use viscosity measurements to discover certain aspects of polymer structure. In the second part you learn the format of a complete laboratory report as well. Before coming to the laboratory, be sure to read pages 106 to 115 in your text book as well as pages 3-41 in your text. The latter reading assignment presents a discussion of error analysis.

Figure 7 – A graph of pV as a function of T



14.2.2 Theory

When the instructors for this course were young (we won't tell you how long ago) there was a TV commercial for Prell Shampoo. To show how "wonderful" Prell was, the advertisement showed two beakers, one filled with Prell and the other filled with water. A pearl was placed at the top of each solution and the rate at which the pearl reached the bottom of the beakers was demonstrated. The pearl took much longer to reach the bottom in the beaker containing Prell than in the beaker containing water.

The difference in rate (which for some reason was supposed to show how good Prell is) is a measure of what we call *viscosity*. Another way of observing the effect of differing viscosity is to pass solutions through a pipe or a narrow capillary. The more viscous the solution the longer it will take for the liquid to pass through the tube. This difference in flow rate is the basic principle behind the Ostwald viscometer that we shall use in this experiment.

The physical origin of viscosity can be understood by imagining two trains of infinite length running on parallel frictionless tracks. Suppose the train on the right track is moving at a higher speed than the train on the left track. Suppose the doors of the trains are open and the passengers on each train participate in a game where they hop back and forth between the two trains at some rate. The passengers jump so that on average the number of people on either train is constant. A person jumping from the fast train to the slow train transfers momentum from the fast train to the slow train. A person jumping from the slow train to the fast train transfers negative momentum. This process of momentum transfer tends to slow the fast train and speed up the slow train. To an observer too far from the train to see the jumping, the change in speed appears to be a frictional drag between the two trains.

Now consider a liquid flowing down a pipe. Imagine the fluid flowing down the pipe to consist of concentric cylindrical layers of decreasing diameter starting at the wall of the pipe and ending at the axis of the pipe. It turns out that the layer adjacent to the wall of the pipe has zero speed and the speed increases from the wall to a maximum at the center of the pipe. Each layer is analogous to the trains described in the preceding paragraph and the molecules that jump between the layers (because of Brownian motion) are analogous to the people jumping between the trains. When we measure the properties of the fluid as it flows down the pipe, we observe a frictional drag between the layers. This frictional drag is the viscosity.

The detailed theoretical background for the experiments is given in the reading assignments. It is useful to note that the key equation for the experiment is given as Eq. (3) in the laboratory text

$$\frac{\eta}{\rho} = Bt \quad (37)$$

where η is the viscosity, ρ is the solution density, B is a constant of proportionality and t is the time necessary for a solution to pass through an Ostwald viscometer.

14.2.3 Experimental Procedure

In this first experiment you limit the measurements to determine the viscosity of the single solution of polyvinyl alcohol by comparison with the viscosity of water. Set up the experimental apparatus as in Figure 1 on page 320 of your text books. Measure the time it takes for water to pass between the marks of the viscometer. You should repeat the measurements for water 10 times. Then perform 10 measurements of the flow time for the polyvinyl alcohol stock solution. Clean the Ostwald viscosimeter thoroughly with 3 portions of distilled water after your polyvinyl alcohol runs. Use a rubber bulb to force water through the capillary tube several times to remove polyvinyl alcohol from the capillary. After the cleaning, repeat the viscosity measurement of distilled water twice. Record the water run time in your data sheet before you ask your TA to sign your data sheet. Fill the Ostwald viscosimeter with distilled water at the end of the laboratory period.

14.2.4 Laboratory Report

In your report you have only three sections; a calculations section, an error analysis section, and a section summarizing your results. Tables of the density and viscosity of water needed for your laboratory reports are found on the last page of this syllabus.

1. Include a cover page giving the title of the experiment and the date the experiment was performed.
2. Include the signed outline of the experimental procedure.
3. Calculations
 - (a) Make a table consisting of the ten values of the flow time for the water measurements and the polyvinyl alcohol solution measurements.
 - (b) Calculate the average flow time for each species.
 - (c) Determine the density and viscosity of water using the tables at the end of this syllabus. As an example, the density of water at 23.1C is 0.997514 g cm⁻³. The viscosity of water at 23.C is 9.38 mP.
 - (d) Use the mean value for water and the known viscosity of water to determine the apparatus constant B .
 - (e) Using the determined value of B , calculate the viscosity of the polyvinyl alcohol solution.
4. Error Analysis
 - (a) Calculate the standard deviation of the mean of the flow time for water.
 - (b) Calculate the standard deviation of the mean of the flow time for the polyvinyl alcohol solution.

- (c) Propagate errors to determine the error in B .
 - (d) Propagate errors to determine the error in the viscosity of the polyvinyl alcohol solution.
5. Summarize your measured results including the calculated errors. Be sure to include the proper number of significant figures and the proper units.
 6. Remember to attach your original data signed by your TA.

Because you need additional experience in error propagation before attempting the next laboratory experiment, do the following exercises as the final section of your laboratory reports:

1. A student has a sodium chloride solution having concentration 1.03 ± 0.02 molar. 50.0 ± 0.1 ml of the solution are placed in a volumetric flask of volume 100.0 ± 0.2 ml, and the solution is mixed with water to the mark on the flask. Determine the final concentration along with its associated random error.
2. The final solution from problem 1 is again diluted in half with water using a pipette and volumetric flask with the same specifications. Determine this final concentration with the associated random error.
3. Later in the semester, you will determine a quantity γ using pressure measurements having the form

$$\gamma = \frac{\frac{P_2}{P_1} - 1}{\frac{P_3}{P_1} - 1}.$$

Suppose the measured pressures are $P_1 = 32.0 \pm 0.2$ torr, $P_2 = 43.2 \pm 0.3$ torr and $P_3 = 762. \pm 1$ torr. Calculate γ and the associated random error. Be careful. In this problem, taking the total differential is essential.

4. In the same experiment later this semester, you will determine a related expression for γ given by

$$\gamma = \frac{\ln \frac{P_2}{P_1}}{\ln \frac{P_3}{P_1}}.$$

Using the same data as in the previous problem, determine γ and the associated random error.

14.3 Experiment 9(28)b: The Intrinsic Viscosity of Polyvinyl Alcohol

14.3.1 Introduction

In this follow-up experiment, you use the viscosity measurements to determine the fraction of polymer units that are connected in the so-called “head to head” arrangement. The approach is to use a reagent that cleaves the head to head connections exclusively. The cleavage increases the number of molecules in solution and decreases the average molecular weight. This change is related to the viscosity. The experimental procedure that you use is identical to that given in the text book. The next subsection will contain a discussion of some of the theoretical ideas.

14.3.2 Theory

We define Δ to be the fraction of bonds connecting monomer units in a strand of polyvinyl alcohol that are connected “head to head.” Let W be the total mass of polyvinyl alcohol in the solution. The total number of monomer units, N_0 , is equal to W divided by M_0 where M_0 is the molecular weight of a monomer; i.e.

$$N_0 = W/M_0 \quad (38)$$

Let N_n be the total number of polymer molecules in the solution before cleavage and N'_n be the total number of polymer molecules in solution after cleavage. If M_n is the molecular weight of the polymer and M'_n is the polymer molecular weight after cleavage, then

$$N_n = W/M_n \quad (39)$$

and

$$N'_n = W/M'_n. \quad (40)$$

Assuming the molecular weights of the polymers are high, to a vanishingly small error we can then write

$$\Delta = \frac{N'_n - N_n}{N_0} \quad (41)$$

$$= \frac{1/M'_n - 1/M_n}{1/M_0} \quad (42)$$

We can then determine Δ from a knowledge of the molecular weights.

We determine the molecular weights using changes in the viscosity that accompany the cleavage process. In a polymer the viscosity increases with increasing molecular weight. The principal reason is that the average volume of the polymer molecules in solution increases with chain length. The quantitative dependence on the molecular weight can be understood by viewing the polymer in a spherical model as discussed in the text book. The final result is found in Eq. (7) on page 323. Since the solution consists of a statistical ensemble, the

molecular weights used are averages resulting in additional correction factors as given in Eq. (11) on page 324 of the text.

Please note that the second of the two equations listed as Eq. (7) on page 323 of the text is in error. You should use

$$M = 7.4 \times 10^4 [\eta]^{1.32}.$$

14.3.3 Cleaning Procedure

As with the first part of this experiment, clean the Ostwald viscosimeter thoroughly with 3 portions of distilled water after your polyvinyl alcohol runs. Use a rubber bulb to force water through the capillary tube several times to remove polyvinyl alcohol from the capillary. After the cleaning, repeat the viscosity measurement of distilled water twice. Record the water run time in your data sheet before you ask your TA to sign your data sheet. Fill the Ostwald viscosimeter with distilled water at the end of the laboratory period.

14.3.4 Laboratory Report

You write complete laboratory reports for this experiment. Before writing your reports, be sure to read pages 10 to 25 in your text books. Your reports should contain the following sections:

1. Title Page: Give the title of the experiment and the date the experiment was performed.
2. Abstract: This should be a 1 paragraph summary of what is to follow including the results. It is suggested that you write the abstract after the rest of the report is complete.
3. Introduction: This should be a discussion of the purpose of the experiment.
4. Theory: This should provide a derivation of the key formulas used in the experiment and used in the analysis of the data. Make sure that all symbols used in this section are defined in words.
5. Procedure: Attach the outline of the procedure you brought to class and had signed by your TA.
6. Original Data: Attach the original data signed by your TA.
7. Results: Tabulate your data so that it can be used in the calculations to follow.
8. Calculations: Provide a sample calculation taking one from the original data to the final result.
9. Error Analysis: Use the standard deviations computed in last weeks experiment as an estimate of the error in the flow times. Determine the error in the specific viscosity by propagation.

10. Graphical Analysis: Plot the specific viscosity divided by the concentration as a function of concentration. Draw the best line through the data. Include error bars on the data points. Use the limiting slope method to determine the intercept and the error in the intercept.
11. Final Calculation: Determine the best average molecular weights for both the cleaved and uncleaved samples. Then calculate Δ . It is possible to determine the error in Δ and the average molecular weight by error propagation. However, we do not expect you to determine these errors in this laboratory report.
12. Summary of Data: Summarize the final results. In particular include the average intrinsic viscosity, the average molecular weight and the best estimate of Δ . Be sure to include the errors for the intrinsic viscosity when reporting your final results, and be sure to report only the proper number of significant figures.
13. Conclusions: Discuss the significance of your results. Be sure to discuss the frequency of the head to head connections and discuss the principal sources of error. Make suggestions for how the experiment might be improved.

15 Format of Laboratory Reports for Final Four Experiments

The laboratory reports for the final four experiments should all have the same format. Each report should have the following sections:

1. Title Page: Give the title of the experiment, the date the experiment was performed and the name of your laboratory partner (if any).
2. Abstract: This should be a 1 paragraph summary of what is to follow including the results. It is suggested that you write the abstract after the rest of the report is complete.
3. Introduction: This should be a discussion of the purpose of the experiment.
4. Theory: This should provide a derivation of the key formulas used in the experiment and in the analysis of the data. Definitions must be provided for all symbols used in this and subsequent sections.
5. Procedure: Attach the outline of the procedure you brought to class and had signed by your TA.
6. Original Data: Attach the original data signed by your TA.
7. Data Table: Tabulate your data so that it can be used in the calculations to follow.

8. Calculations: Provide a sample calculation taking one from the original data to the final results.
9. Graphical Analysis: If required, plot your data on a graph. Draw the best line through the data. Determine the slope and/or intercept of the graph as needed.
10. Error Analysis: Perform any necessary propagation of errors. You need show only one example of any one kind of propagation.
11. If a graph is needed use the calculated errors to include error bars on the graph. Use the limiting slope method to determine the error in the slope and/or intercept.
12. Propagate the errors from the graph to give the errors in the final results.
13. Summary of Data: Summarize the final result with error bars included. Be sure to report only the proper number of significant figures.
14. Conclusions: Discuss the significance of your results. Make suggestions for how the experiment might be improved.

Special requirements for the last four experiments are summarized in the next section.

16 Second Four Experiments

16.1 Experiment Number 3(3) - Heat Capacity Ratio of Gases

16.1.1 Reading Assignment

Read experiment number 3(3), pages 51-59 in the text. The discussion that follows is somewhat different and more complete than the discussion in the textbook.

16.1.2 Theoretical Discussion

To determine the ratio $\gamma = C_p/C_V$, a carboy is filled with some gas at an initial pressure p_1 , and initial temperature T_1 and an initial volume V_1 . The initial pressure is generally about 60mm in excess of the external atmospheric pressure. After an equilibration period where the initial pressure is found to be unchanged for several minutes, the stopper at the top of the carboy is removed for about 1 second and replaced. It is assumed the rapidity that the stopper is removed and replaced allows the gas to expand adiabatically against the external atmospheric pressure p_2 to an intermediate volume V_2 and a temperature T_2 . Because the expansion is adiabatic, we know that $T_2 < T_1$. Finally, the gas heats at constant volume to the initial temperature T_1 so that the final pressure is p_3 . The experimental data consist of the pressures p_1 , p_2 and p_3 .

- Assumption of Reversibility: According to the textbook, the adiabatic expansion can be viewed as reversible. Assuming reversibility the determination of γ follows from the standard expressions for the adiabatic expansion of an ideal gas. Because the expansion is adiabatic and assumed reversible, we have

$$dU = -pdV = -\frac{nRT}{V}dV. \quad (43)$$

For an ideal gas

$$dU = C_V dT \quad (44)$$

so that

$$-nR d \ln V = C_V d \ln T. \quad (45)$$

Then for the first step in the process

$$-nR \ln \frac{V_2}{V_1} = C_V \ln \frac{T_2}{T_1}. \quad (46)$$

Using the ideal gas law we have

$$\frac{p_2 V_2}{T_2} = \frac{p_1 V_1}{T_1} \quad (47)$$

so that

$$\frac{T_2}{T_1} = \frac{p_2 V_2}{p_1 V_1}. \quad (48)$$

Then

$$-nR \ln \frac{V_2}{V_1} = C_V \ln \frac{p_2 V_2}{p_1 V_1} = C_V \ln \frac{p_2}{p_1} + C_V \ln \frac{V_2}{V_1}. \quad (49)$$

Using the ideal gas relation

$$C_p = C_V + nR \quad (50)$$

we then obtain

$$-C_p \ln \frac{V_2}{V_1} = C_V \ln \frac{p_2}{p_1}. \quad (51)$$

In the second step of the process we heat the gas at constant volume to a final state defined by the variables p_3 , V_2 and T_1 . For this constant volume process the ideal gas law along with Eq. (48) give

$$\frac{T_2}{T_1} = \frac{p_2 V_2}{p_1 V_1} = \frac{p_2}{p_3} \quad (52)$$

so that

$$\frac{V_2}{V_1} = \left(\frac{p_1}{p_2} \right) \left(\frac{p_2}{p_3} \right) = \frac{p_1}{p_3}. \quad (53)$$

Then

$$-C_p \ln \frac{p_1}{p_3} = C_V \ln \frac{p_2}{p_1} \quad (54)$$

or

$$\gamma = \frac{C_p}{C_V} = \frac{\ln(p_1/p_2)}{\ln(p_1/p_3)}. \quad (55)$$

- Irreversible Adiabatic Expansion: We now (perhaps more realistically) assume that the adiabatic expansion is an irreversible expansion of the gas against the constant external atmospheric pressure p_2 . Equating the irreversible work to the energy change for an ideal gas, we obtain

$$C_V(T_2 - T_1) = -p_2(V_2 - V_1) \quad (56)$$

$$= -p_2 \left(\frac{nRT_2}{p_2} - \frac{nRT_1}{p_1} \right) \quad (57)$$

$$= -nR \left(T_2 - T_1 \frac{p_2}{p_1} \right) \quad (58)$$

or

$$C_V \left(\frac{T_2}{T_1} - 1 \right) = -nR \left(\frac{T_2}{T_1} - \frac{p_2}{p_1} \right). \quad (59)$$

Using Eq. (48)

$$C_V \left(\frac{p_2 V_2}{p_1 V_1} - 1 \right) = -nR \left(\frac{p_2 V_2}{p_1 V_1} - \frac{p_2}{p_1} \right). \quad (60)$$

The second step in the irreversible process is the same as the reversible case, and we can again use Eq.(53) to obtain

$$C_V \left(\frac{p_2 p_1}{p_1 p_3} - 1 \right) = -nR \left(\frac{p_2 p_1}{p_1 p_3} - \frac{p_2}{p_1} \right) \quad (61)$$

or from Eq.(50)

$$C_V \left(\frac{p_2}{p_3} - 1 \right) = -(C_p - C_V) \left(\frac{p_2}{p_3} - \frac{p_2}{p_1} \right) \quad (62)$$

or

$$-(\gamma - 1) \left(\frac{p_2}{p_3} - \frac{p_2}{p_1} \right) = \left(\frac{p_2}{p_3} - 1 \right). \quad (63)$$

After rearranging Eq. (63) and solving for γ we obtain the final result

$$\gamma = \frac{p_1/p_2 - 1}{p_1/p_3 - 1}. \quad (64)$$

We can obtain Eq. (64) by expanding the numerator and denominator of Eq.(55) for small values of the argument. If we write

$$f(x) = \ln x \quad (65)$$

and expand f in a Taylor series about $x = 1$, we obtain

$$f(x) = f(1) + (x - 1)f'(1) + \frac{1}{2}(x - 1)^2 f''(1) + \dots \quad (66)$$

or

$$\ln x = (x - 1) + \dots \quad (67)$$

Then for $(p_1/p_2 - 1) < 1$ and $(p_1/p_3 - 1) < 1$ to leading order Eq.(55) becomes Eq.(64).

- Correct derivation allowing the escape of some of the gas: Let us now carefully derive Eq.(64). The steps in the process are:

1. an adiabatic irreversible expansion of the ideal gas from the state n_1, p_1, T_1 and V_1 against a constant external pressure p_2 to the final state n_1, p_2, T_2 and V_2 ;
2. use of the initial volume of the container so that the state of the system is represented by n_2, p_2, V_1 and T_2 - in this step some of the gas originally present is allowed to escape into the atmosphere;
3. constant volume heating to the initial temperature represented by the final state n_2, V_1, p_3 and T_1 .

We let C_V and C_p be the constant volume and constant pressure heat capacities *per mole*. Then for step 1

$$n_1 C_V (T_2 - T_1) = -p_2 \left(\frac{n_1 R T_2}{p_2} - \frac{n_1 R T_1}{p_1} \right) \quad (68)$$

$$= -n_1 R T_2 + n_1 R T_1 \frac{p_2}{p_1}. \quad (69)$$

Solving for T_2 we obtain

$$n_1 (R + C_V) T_2 = n_1 R T_1 \frac{p_2}{p_1} + n_1 C_V T_1 \quad (70)$$

or

$$C_p T_2 = R T_1 \frac{p_2}{p_1} + C_V T_1 \quad (71)$$

so that

$$T_2 = \frac{R}{C_p} T_1 \frac{p_2}{p_1} + \frac{C_V}{C_p} T_1 \quad (72)$$

$$= \frac{C_p - C_v}{C_p} T_1 \frac{p_2}{p_1} + \frac{C_V}{C_p} T_1 \quad (73)$$

$$= \left(1 - \frac{1}{\gamma} \right) T_1 \frac{p_2}{p_1} + \frac{1}{\gamma} T_1 \quad (74)$$

$$= \frac{T_1}{\gamma} \left[(\gamma - 1) \frac{p_2}{p_1} + 1 \right]. \quad (75)$$

In the second step, we remove $n_1 - n_2$ moles of gas from our system. The expression for n_2 is

$$n_2 = \frac{p_2 V_1}{R T_2} \quad (76)$$

$$= \frac{p_2 V_1}{R} \frac{\gamma}{T_1 \left[(\gamma - 1) \frac{p_2}{p_1} + 1 \right]}. \quad (77)$$

In step 3, the gas remaining in the initial volume is heated to temperature T_1 so that its final pressure is p_3 . Solving for p_3

$$p_3 = \frac{n_2 RT_1}{V_1} \quad (78)$$

$$= \left(\frac{RT_1}{V_1} \right) \left(\frac{p_2 V_1}{R} \right) \left(\gamma / \left\{ T_1 \left[(\gamma - 1) \frac{p_2}{p_1} + 1 \right] \right\} \right) \quad (79)$$

or

$$p_3 = \frac{p_2 \gamma}{(\gamma - 1) \frac{p_2}{p_1} + 1}. \quad (80)$$

We next rearrange Eq.(80)

$$(\gamma - 1) \frac{p_2}{p_1} + 1 = \gamma \frac{p_2}{p_3} \quad (81)$$

and solve for γ

$$\gamma \left(\frac{p_2}{p_1} - \frac{p_2}{p_3} \right) = \frac{p_2}{p_1} - 1 \quad (82)$$

or

$$\gamma = \frac{p_1/p_2 - 1}{p_1/p_3 - 1} \quad (83)$$

in complete agreement with Eq.(64).

16.1.3 Experimental Notes

1. The adiabatic expansion method is used. Follow the procedure outlined in the text.
2. Our experimental design uses an electronic open ended manometer. Readings from an open ended manometer gives the pressure relative to the atmospheric pressure. The atmospheric pressure must be obtained by reading the electronic barometer. Ask your TA for instructions.

16.1.4 Laboratory Report

1. Theory Section:
 - (a) Be sure to discuss the difference between the assumption of a reversible adiabatic expansion and an irreversible expansion against a constant external pressure. Explain which assumption is closest to the actual experimental situation.
2. Calculations:
 - (a) Calculate γ using both Eqs. (55) and (64).
3. Error Analysis:

- (a) Assume an error of ± 2 in the last recorded figure from the electronic manometer and the electronic barometer.
- (b) Although the reported values of γ are obtained by an average of the heat capacity ratio for three trials, choose only 1 of the 3 trials for each gas to propagate the error in γ . Make sure to indicate carefully which of the three trials (for each gas) is used in your error propagation calculation. Carefully propagate errors in the pressure, the pressure ratios, the logarithms and the final ratios to obtain the propagated error in the heat capacity ratio. Compute expressions and values for the statistical error in γ from both Eqs. (55) and (64).

4. Conclusions:

- (a) Compare $\gamma \pm \epsilon(\gamma)$ using both the assumption of reversibility and irreversible adiabatic expansions. Do the two values determined lie within their respective statistical errors?
- (b) Compare $\gamma \pm \epsilon(\gamma)$ with the theoretical values expected for an ideal monatomic and diatomic gas. Do the theoretical results lie within your calculated experimental errors? If the results lie outside the experimental errors, discuss possible reasons for the lack of agreement.

16.2 The Phase Transitions in Ammonium Chloride

16.2.1 Reading Assignment

Engel and Reid, Chapters 7 and 8

16.2.2 Introduction

Common, every-day physical phenomena like the melting of an ice cube or the boiling of water are examples of phase transitions. While such phase transitions are ubiquitous and important, in addition to melting, boiling and sublimation, there are other varieties of phase transitions that have interesting phenomenology worthy of study. This experiment examines the richness of the kinds of phase transitions using temperature dependent studies of ammonium chloride crystals. At atmospheric pressure ammonium chloride exhibits both a sublimation transition and additionally a different category of phase transition often called an order to disorder transition. In the context of diverse phase transition phenomena, ammonium chloride represents a useful and illustrative test case.

All phase transitions are characterized by discontinuities or infinities in certain physical properties. For example when a solid melts there is a visible phase boundary between the solid and liquid phases. To be concrete an ice cube floating in water at 0°C and 1 atmosphere pressure exhibits a visible boundary. The boundary between the ice and water is visible, because the index of refraction of ice and water are different even though both phases are

at the same temperature and pressure. The index of refraction is discontinuous across the phase boundary separating the solid and liquid phases rendering a visible surface. Other properties that are discontinuous across the phase boundary are the molar volume, the density, the molar constant pressure heat capacity, and the molar entropy. Not only is C_P discontinuous across the phase boundary (i.e. the molar heat capacities of ice and water differ), there is also an infinity in C_P at the transition temperature. The process of melting a solid is always isothermal with an associated latent heat $\Delta H_{m, fus}^\circ$ (the enthalpy of fusion), and $C_P = (\partial H/\partial T)_P$ must be infinite owing to the finite latent heat with no associated temperature change.

16.2.3 Orders of Phase Transitions

In addition to phase transitions like the melting of a solid or the boiling of a liquid, there are other kinds of phase transitions in nature where there is no latent heat or discontinuity in the density. These other kinds of phase transitions are associated with infinities or discontinuities in properties other than the density or molar entropy. To understand these alternate kinds of behaviors, Paul Ehrenfest has introduced a classification scheme for phase transitions based on derivatives of the chemical potential of a substance. We can label two phases that are in equilibrium with the Greek letters α and β . For example α might represent a liquid phase and β a solid phase. Because at equilibrium any two systems of the same chemical composition must have the same chemical potential μ

$$d\mu_\alpha = d\mu_\beta. \quad (84)$$

Recalling that for a pure chemical substance the chemical potential is the molar Gibbs free energy and recalling the total differential of the Gibbs free energy

$$d\mu = -S_m dT + V_m dP, \quad (85)$$

Eq. (84) implies that

$$-S_{m,\alpha} dT + V_{m,\alpha} dP = -S_{m,\beta} dT + V_{m,\beta} dP. \quad (86)$$

In many, common phase transitions $S_{m,\alpha} \neq S_{m,\beta}$ and $V_{m,\alpha} \neq V_{m,\beta}$. For example the molar entropy of liquid water and ice are not the same, and densities of ice and water differ. Both the entropy and molar volume of a substance can be computed from the first derivatives of the chemical potential

$$S_m = \left(\frac{\partial \mu}{\partial T} \right)_V \quad (87)$$

and

$$V_m = \left(\frac{\partial \mu}{\partial P} \right)_T. \quad (88)$$

When the first derivatives of the chemical potential are discontinuous across a phase boundary, the phase transition is called a *first-order phase transition*. Across a phase boundary, such first-order phase transitions have a discontinuous density and molar entropy

$$\Delta S_m = S_{m,\beta} - S_{m,\alpha} = \frac{\Delta H_{m,trans}}{T_{trans}} \quad (89)$$

where T_{trans} is the transition temperature and $\Delta H_{m,trans}$ is the latent molar heat (enthalpy change) for the transition. The melting of ice and the boiling of liquid water are examples of first-order transitions, because both transitions have a latent heat and a density discontinuity across the phase boundary.

We can imagine phase transitions that have continuous first derivatives of the chemical potential but discontinuous second derivatives. When the first derivatives of the chemical potential are continuous but the second derivatives are discontinuous, the transition is called a *second-order phase transition*. In second-order phase transitions, the discontinuities are associated with the following three thermodynamic properties:

$$\left(\frac{\partial^2 \mu}{\partial T^2} \right)_P = - \left(\frac{\partial S_m}{\partial T} \right)_P \quad (90)$$

$$= \frac{C_{m,P}}{T} \quad (91)$$

where $C_{m,P}$ is the molar heat capacity at constant pressure,

$$\left(\frac{\partial^2 \mu}{\partial P^2} \right)_T = \left(\frac{\partial V_m}{\partial P} \right)_T \quad (92)$$

$$= -V_m \kappa \quad (93)$$

where κ is the isothermal compressibility, and

$$\left(\frac{\partial}{\partial T} \left(\frac{\partial \mu}{\partial P} \right)_T \right)_P = \left(\frac{\partial V_m}{\partial T} \right)_P \quad (94)$$

$$= V_m \beta_V \quad (95)$$

where β_V is the isobaric coefficient of thermal expansion.

As of this writing the only true second-order transition that is believed to have been observed in nature is the normal to superconducting transition in metals at zero magnetic field. This transition from normal to superconducting behavior has no latent heat, no density discontinuity, but the constant pressure heat capacity, the isobaric coefficient of thermal expansion and the isothermal compressibility are all discontinuous.

There is a class of phase transitions that is neither first-order nor second-order and is associated with slow-growing infinities in certain thermodynamic properties. Such phase transitions, often correctly called *continuous phase transitions* or λ *transitions* (a graph of

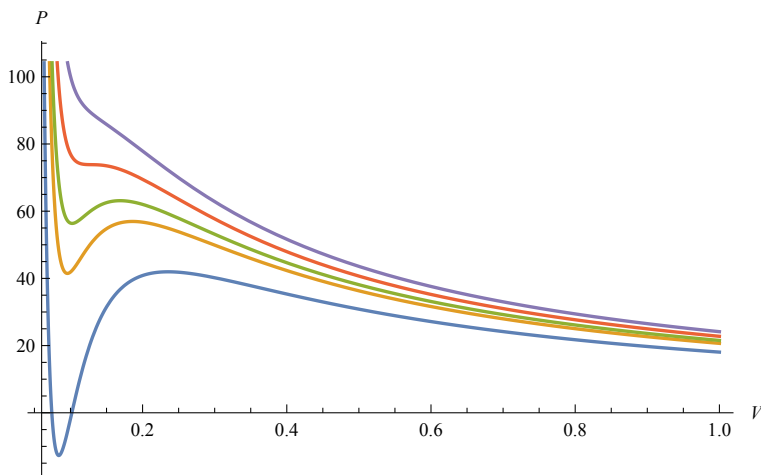


Figure 8: Van der Waals isotherms for CO_2 . The curves from bottom to top correspond to $T = 250. \text{ K}$, $T = 280. \text{ K}$, $T = 290. \text{ K}$, $T = 304.12 \text{ K}$ and $T = 320. \text{ K}$. The critical temperature is 304.12 K

C_P as a function of T near the transition temperature has the shape of the Greek letter λ), are sometimes incorrectly referred to as second-order transitions. In the sense of Ehrenfest, continuous transitions are not really second order. Although many continuous phase transitions are associated with order to disorder phenomena, these continuous phase transitions are formally identical to the behavior of fluids at critical points. In statistical mechanics courses it is often proved that singularities in some thermodynamic properties of fluids at critical points match the behavior of the singularities in systems undergoing an order to disorder transition. We will accept the connection and use the singularities at the critical point of fluids to introduce the kinds of singularities expected in the order to disorder transition in ammonium chloride.

16.2.4 Critical Phenomena in the Van der Waals Fluid

The Van der Waals equation of state

$$P = \frac{RT}{V_m - b} - \frac{a}{V_m^2} \quad (96)$$

is a pedagogically useful approximate equation of state. Generally, P , V and T data are fit for specific gases to determine the two numerical parameters a and b . The b parameter approximately corrects the ideal gas law for the finite volume occupied by the constituent molecules of the gas, and the a parameter approximately corrects for the nonzero intermolecular forces between the molecules. Figure 8 depicts typical Van der Waals isotherms for $a=3.658 \text{ bar mol}^{-2} \text{ L}^2$ and $b=0.0429 \text{ L mol}^{-1}$, where the parameters are chosen to fit data for carbon dioxide.

The Van der Waals equation of state is an example of a *mean field* theory. The term “mean field” implies that the corrections for non ideality are included in an average sense by having each particle interact with the average density distribution of all other particles.

Unlike the ideal gas law, which applies only to gases at low pressures, the Van der Waals equation applies to both the liquid and vapor phases. In Figure 8 the lowest temperature isotherm presents a highly incompressible phase at high pressures [i.e. $(\partial P/\partial V)_T$ is very large] and a highly compressible phase at low pressures [i.e. $(\partial P/\partial V)_T$ is small]. The incompressible phase is the liquid phase, and the compressible phase is vapor. Between the liquid and vapor phases is a region where

$$\left(\frac{\partial P}{\partial V_m}\right)_T > 0, \quad (97)$$

which is impossible. It is usual to interpret the unphysical region where the volume and pressure increase together as the liquid-vapor coexistence region; i.e. the location of the liquid to vapor phase transition. Because the region where $(\partial P/\partial V)_T > 0$ is unphysical, it is common to replace the so-called Van der Waals loops with a horizontal straight line as depicted in Fig. 9. The resulting isotherm is then qualitatively the same as a real isotherm where the distance along the V axis of the horizontal line is identical to the discontinuous change in the molar volume between the gas and liquid phases. To locate the horizontal line in the case of a Van der Waals fluid, it is usual to use the Maxwell construction where

$$\Delta G_{m,vap} = \int_{V_{m,\ell}}^{V_{m,g}} dG = \int_{V_{m,\ell}}^{V_{m,g}} V_m dP \quad (98)$$

is set to 0; i.e. the condition for equilibrium between the phases. In Eq. (98) $V_{m,\ell}$ is the molar volume of the liquid and $V_{m,g}$ is the molar volume of the gas phase. Using integration by parts it can be proved that $\Delta G_{m,vap} = 0$ when the area under the Van der Waals loop matches the area under the horizontal line.

Real isotherms for fluids must contain discontinuities as seen in Fig. 9 when the loop region is replaced by the horizontal line included using the Maxwell construction. Consequently, in contrast to the ideal gas law, which is valid for all gases in the limit of zero pressure, the Van der Waals equation of state is valid for no real substance. However, much can be learned about real fluids when the Van der Waals isotherms are properly interpreted.

By examining Fig. 8 we see that the loop region becomes increasingly narrow as the temperature is increased. Using the Maxwell construction $\Delta V_{m,vap}$ then decreases with increasing temperature. In real fluids $\Delta V_{m,vap}$ also decreases with temperature until a temperature is reached at which $\Delta V_{m,vap}=0$. The temperature at which $\Delta V_{m,vap}=0$ is called the *critical temperature* often defined as the temperature above which liquification cannot occur.

By examining Fig. 8 it is evident that the loop region disappears when the maximum and minimum in the loop region coalesce into a single point. Such a point is called a point of inflection, and the conditions for the inflection point are the two requirements $(\partial P/\partial V_m)_T = (\partial^2 P/\partial V_m^2)_T = 0$. By differentiating Eq. (96) twice to satisfy the conditions

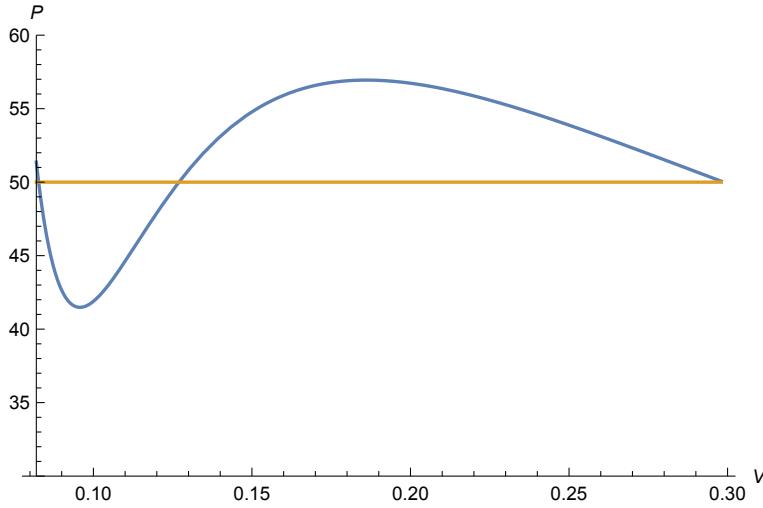


Figure 9: The Van der Waals isotherm at 280 K showing the Maxwell construction for the liquid to vapor coexistence region

for the critical point, the critical constants for the Van der Waals fluid can be found

$$V_c = 3b \quad T_c = \frac{8a}{27Rb} \quad P_c = \frac{a}{27b^2} \quad (99)$$

Because $(\partial P/\partial V_m)_T$ vanishes at the critical point, the isothermal compressibility must diverge at the critical temperature; i.e.

$$\lim_{T \rightarrow T_c} \kappa = \lim_{T \rightarrow T_c} \left[-\frac{1}{V_m} \left(\frac{\partial V_m}{\partial P} \right)_T \right] = \infty \quad (100)$$

An important question is how rapidly κ diverges in the limit along the critical isochore (where V is constant and equal to the critical volume). Evaluating the first derivative of Eq. (96) at $V_c = 3b$ we have

$$\left(\frac{\partial P}{\partial V_c} \right)_T = -\frac{RT}{4b^2} + \frac{2a}{27b^3} \quad (101)$$

Letting

$$T = T_c + \epsilon \quad (102)$$

where

$$\epsilon = T - T_c \quad (103)$$

and substituting the expression for T into Eq. (101)

$$\lim_{T \rightarrow T_c} \left(\frac{\partial P}{\partial V_m} \right)_T = -\frac{R}{4b^2} \epsilon \quad (104)$$

or

$$\lim_{T \rightarrow T_c} \kappa = A|T - T_c|^{-\gamma} \quad (105)$$

where for a Van der Waals fluid $\gamma = 1$. Equation (105) is an example of a *power law*; i.e. a property that depends on another property to some power. The exponent γ is an example of a *critical exponent*. Critical exponents express the divergence of some property as the critical point is approached. For the Van der Waals fluid along the critical isochore the isothermal compressibility diverges inversely as $|T - T_c|$ to the first power. As we have implied other thermodynamic properties have discontinuities that obey power laws in the vicinity of a critical point. For example, along the critical isochore the constant pressure heat capacity diverges as

$$\lim_{T \rightarrow T_c} C_P = B|T - T_c|^{-\gamma'}. \quad (106)$$

Because the critical behavior of C_P is important to the ammonium chloride experiment, we now show that $\gamma' = 1$ for a Van der Waals fluid along the critical isochore where V is constant and equal to $V_c = 3b$.

We start with the total differential of the Helmholtz free energy

$$dA = -SdT - PdV \quad (107)$$

and write

$$\left(\frac{\partial A}{\partial V}\right)_T = -P, \quad (108)$$

which for the Van der Waals fluid gives

$$\left(\frac{\partial A}{\partial V}\right)_T = -\frac{nRT}{V - nb} + \frac{an^2}{V^2}. \quad (109)$$

An expression for the Helmholtz free energy under isothermal conditions can be obtained by integrating

$$dA = -\frac{nRT}{V - nb}dV + \frac{an^2}{V^2}dV \quad (110)$$

to obtain

$$A = -nRT \ln(V - nb) - \frac{an^2}{V} + C(T) \quad (111)$$

where $C(T)$ is some temperature-dependent constant. The Helmholtz free energy for an ideal gas can be obtained from Eq. (111) by setting a and b to 0

$$A^{ideal} = -nRT \ln V + C(T) \quad (112)$$

so that

$$A - A^{ideal} = -nRT \ln \frac{V - nb}{V} - \frac{an^2}{V} \quad (113)$$

Now using Eq. (107)

$$\left(\frac{\partial A}{\partial T}\right)_V = -S \quad (114)$$

and

$$\left(\frac{\partial S}{\partial T}\right)_V = \frac{C_V}{T}. \quad (115)$$

Equation (113) then gives

$$-\left(\frac{\partial^2(A - A^{ideal})}{\partial T^2}\right)_V = \frac{C_V - C_V^{ideal}}{T} = 0 \quad (116)$$

or the Van der Waals constant volume heat capacity matches the ideal gas result, which is free of divergences. However, we now show that C_P does diverge at the critical point. Using

$$C_P - C_V = T \left(\frac{\partial V}{\partial T}\right)_P \left(\frac{\partial P}{\partial T}\right)_V \quad (117)$$

the Van der Waals equation of state gives

$$C_P - C_V = \frac{nR}{1 - \frac{2an(V - nb)^2}{RTV^3}}. \quad (118)$$

Along the critical isochore we set $V = V_c = 3b$ in Eq. (118) and obtain

$$\frac{C_P - C_V}{nR} = \frac{1}{1 - \frac{8a}{27RbT}} \quad (119)$$

$$= \frac{T}{T - \frac{8a}{27Rb}} \quad (120)$$

$$\equiv \frac{T}{T - T_c} \quad (121)$$

Because C_V has no divergences, C_P diverges as

$$\lim_{T \rightarrow T_c} C_p = B|T - T_c|^{-\gamma'} \quad (122)$$

with $\gamma' = 1$ along the critical isochore.

The use of the “critical isochore” requires some explanation. The values of the critical exponents depend on the choice of constraint used in their evaluation. If the critical exponent were evaluated along the critical isobar where P is constant and set such that $P = P_c = a/27b^2$, for the Van der Waals fluid the critical exponent for C_P is $\gamma' = 2/3$. Owing to the cubic dependence of the volume on P and T , the demonstration of the critical exponent of C_P for the Van der Waals fluid is more complex than the preceding discussion, and the details are found in the Appendix.

When critical exponents are determined experimentally for real fluids, the measured exponents do not match those calculated for the Van der Waals fluid. Evidently, the mean-field approximations that lead to the Van der Waals equation of state are too severe to capture the true behavior at the critical point. For example, along the critical isochore the compressibility for a real fluid diverges with an exponent $\gamma = 1.23 - 1.25$ and the critical exponent for C_P is the same; i.e. $\gamma' = 1.23-1.25$. In fact, the result $\gamma = \gamma'$ is quite general. Critical exponents are interrelated, and while there are many diverging thermodynamic properties at the critical point, there are fewer independent exponents. The interrelations between the exponents at the critical point arise from something called *scaling*.

Experimentally and theoretically it is understood that all real fluids have the same critical exponents. In other words, the measured values of γ or γ' are the same whether the fluid is argon, carbon dioxide or water. In this sense critical exponents are universal. In fact the same exponents are found for fluids, some magnetic materials and ammonium chloride. Fluids, some magnetic material, ammonium chloride, and other systems that share critical exponents are said to belong to the same *universality class*.

16.2.5 Measuring Critical Exponents

Ammonium chloride undergoes an order to disorder transition at around 242K. The measurement of the critical exponent associated with C_P in ammonium chloride is an important goal of this experiment. Unfortunately, accurate measurements of critical exponents are difficult, and the accurate determination of the exponents is a topic of current research.

There are multiple reasons that determining critical exponents is difficult, and we discuss a few of the reasons now. One problem is the singularity in the property of interest often has a non-singular background contribution. For example for the Van der Waals fluid along the critical isochore and near the critical point

$$C_P = C_V + B|T - T_c|^{-1} \quad (123)$$

the non-singular portion is the constant volume heat capacity. To measure the singular part properly, the non-singular part must be known, or data must be taken so close to the critical point that the non-singular part becomes unimportant. Because the background is usually not known, high quality data are then needed very close to the singular point.

A second, perhaps more serious problem, is the critical temperature is not known *a priori*. Consequently, both the exponent and T_c need to be fit simultaneously to the data. Such two-parameter fitting to high precision can be difficult.

Another serious problem is often called *critical slowing down*. As one chooses temperatures that are closer and closer to the critical point, it takes increasing amounts of time to equilibrate the system. The amount of time, then, to gather data near the critical point can grow making accurate data difficult to determine.

In spite of the difficulties in this experiment we will determine the critical exponent for C_P in the order to disorder transition in NH_4Cl . In evaluating the precision of the results, it is important to be aware of the difficulties associated with any determination of critical

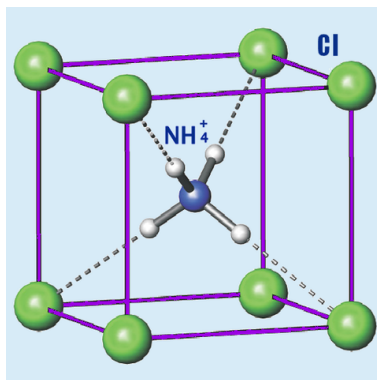


Figure 10: A representation of the NH_4Cl crystal taking on the CsCl lattice type. The hydrogen atoms on the ammonium ion point toward the chloride ions owing to attractive interatomic forces. The figure is taken from Wikipedia.

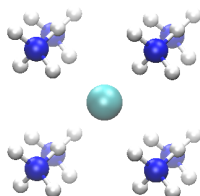


Figure 11: A segment of the ammonium chloride lattice with a chloride ion surrounded by ammonium ions in a cubic arrangement.

exponents. We should not be disappointed if our determined exponent is only approximately correct.

16.2.6 The Order-disorder Transition in NH_4Cl

A representation of the ammonium chloride lattice is given in Fig. 10. Ammonium chloride has a CsCl lattice type with eight, near-neighbor chloride anions occupying the vertex locations of a simple cube with each ammonium ion in the center of the cube. Additionally near neighbor ammonium ions also form a cube with chloride ions in the center of the ammonium chloride cubes as depicted in Fig. 11. The lowest energy orientation of the ammonium ions has each hydrogen atom pointing directly to a chloride ion as seen Fig. 10. Additionally the ammonium ions interact with each other so that at lowest energy all the ammonium

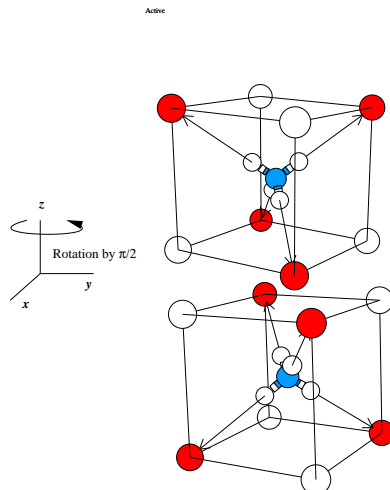


Figure 12: The two crystallographically distinct orientations of the ammonium ions in the NH_4Cl lattice. In the top figure the ammonium ion “points” to the darkened chloride ions, and in the bottom figure the ammonium ions point to the second set of chloride ions in the lattice. Owing to the interactions between the ammonium ions, at 0K, all ammonium ions are oriented to crystallographically identical chloride ions.

ions in the lattice point to crystallographically identical sets of chloride ions. The different crystallographic orientations are displayed in Fig. 12. In this fashion at 0K all ammonium ions are oriented in the same direction and the lattice is completely ordered.

At non-zero temperatures the ammonium ions are not perfectly ordered owing to thermal fluctuations. These fluctuations disrupt the ordering of the lattice so that at the lowest attainable temperatures a few of the ammonium ions are oriented toward the second set of crystallographically identical chloride ions. As the temperature increases, while a majority of NH_4^+ ions still orient toward the same set of chloride ions, increasing numbers of ammonium ions orient toward the other set of chloride ions. At the transition temperature for the order-disorder transition an equal number of ammonium ions orient toward both crystallographically equivalent sets of chloride ions. Above the transition temperature the orientations remain completely random.

The behavior of NH_4Cl is similar to the behavior of some magnetic materials. A simple model (called the Ising model) of magnetic materials considers a simple cubic lattice with spins at each lattice site. Each spin acts like a small bar magnet, and the spins can either be spin up or spin down. Near-neighbor spins are allowed to interact with each other, and for ferromagnetic materials at absolute zero all spins point in the same direction (e.g. all are spin up). At low but non-zero temperatures the thermal fluctuations induce some of the spins to “point” down. As the temperature increases, increasing numbers of spins point down until a transition temperature where the numbers of up and down orientations are

random. The random orientations continue as the temperature is increased past the transition temperature. The comparison between the behavior of the Ising model and ammonium chloride helps explain why NH_4Cl and the Ising system have the same critical exponents; i.e. both systems belong to the same universality class.

16.2.7 Experimental Procedure

The measurements are made with a differential scanning calorimeter (DSC) in the Teknor Apex lab. The DSC is quite delicate, not to mention expensive, and no parts of this procedure should be carried out without full supervision by one of the instructors.

You will need 2 pair of gloves for this experiment. Make sure to purchase gloves from the Stockroom before coming to lab

1. Obtain the sample pans from the instructor. Wearing gloves, weigh out between 5 and 10 mg of ammonium chloride in the sample pan. Make sure to record the precise mass of the ammonium chloride to at least 4 decimal places. The balance will report the mass in grams, but you will need to enter the mass using the DSC software in mg.
2. Place the lid on the sample pan, and crimp the pan using the press as shown by your instructor.
3. Place the sample pan in the DSC as demonstrated by your instructor.
4. Remove your gloves before entering data into the computer connected to the DSC.
5. Enter the measurement parameters into the computer following the instructions on the sheet at the side of the computer. First carry out an initial scan at a rapid rate to locate the temperature of the order to disorder transition using the analysis on the screen. Once the transition temperature is identified, begin the final scan at about 2 degrees Kelvin below the transition temperature.
6. Start the final scan. The total time to complete the scan should be about 2 1/2 hours. You may leave the room while the scan is proceeding.
7. Using the data displayed graphically on the computer screen, identify the low-temperature order-disorder transition and the high-temperature sublimation transition. Use the screen software to determine the enthalpy of sublimation per gram for the high temperature transition. Make sure to record this enthalpy.
8. Your instructor will help you retrieve your electronic data and e-mail the data in ascii format to you.

16.2.8 Calculations

Unlike all other experiments this semester, you are allowed to use any spread sheet software (Excel, Google Sheets, etc.) to perform the calculations and the graphical analysis.

1. Import your ascii file into your spread sheet software. The first column is the time in seconds, the second is the temperature in degrees C and the third is the power in mW(mJs^{-1}).
2. Locate the low-temperature peak in the power (a negative number) and determine the temperature at the peak. This temperature is the critical temperature for the order-disorder transition. In the vicinity of the low-temperature heat capacity peak, the rate of temperature change is a constant. Because

$$C_P = \frac{dq}{dT} = \frac{dq}{dt} \frac{dt}{dT}, \quad (124)$$

where t is the time and T is the temperature, the negative of the power is proportional to the heat capacity. We are to determine the critical exponent, so the absolute magnitude of the heat capacity is not required for the calculations. You can consider the power to be C_P .

3. Use the spread sheet to convert the temperatures to degrees K.
4. Use the spread sheet to calculate $\ln |T - T_c|$ for 20 data points above and 20 data points below the critical temperature T_c .
5. Use the spread sheet to calculate $\ln C_P$ for the 40 data points.
6. Using the spread sheet plot C_P as a function of T in degrees K for the 41 data points from above.
7. Plot $\ln C_P$ as a function of $\ln |T - T_c|$ for the 10 points below the critical temperature. Choose a region of linearity away from the peak to determine the slope of the line that connects these points. The slope is your low-temperature critical exponent.
8. Repeat the determination of the critical exponent on the high-temperature side of the heat capacity maximum.

16.2.9 Error Analysis

There is no required error analysis for this laboratory report.

16.2.10 Theory Section of Laboratory Report

Because there is no error analysis section for this report, a more extensive theoretical section than other laboratory reports is expected. Give as much detail as possible concerning the concept of a critical point using the Van der Waals fluid as a model. In the discussion complete the integration by parts required in Eq. (98) and complete the derivations of Eqs. (99) and (118).

16.2.11 Discussion

Do the critical exponents, γ_ℓ and γ_h associated with the low-temperature and high-temperature sides of the order to disorder heat capacity peak agree (they should)? What does the agreement or lack of agreement tell you about the accuracy of the approach used in this experiment to determine the critical exponents? How close are the values of γ determined experimentally to the values given for the Van der Waals fluid? In answering this last question should the critical exponent along the critical isochore or isobar be used (if either)? Make sure to clarify which exponent is chosen and why in the discussion.

Compare your measured sublimation transition temperature and molar enthalpy of sublimation with literature values.

16.3 Experiment 7(13) - The Vapor Pressure of a Pure Liquid

16.3.1 Reading Assignment

Read experiment number 7(13), pages 85-93 in the text.

16.3.2 Experimental Notes

1. The boiling point method is used. Follow the procedure in the book.
2. Our experimental design uses an open ended electronic manometer. Readings from an open ended manometer gives the pressure relative to the atmospheric pressure. The atmospheric pressure must be obtained by reading the electronic barometer.
3. No baffle is used.
4. We use the internal vacuum line (colored orange in our building) for the vacuum pump.
5. Glassware under vacuum is dangerous if broken. Avoid tampering with the connections when the system is evacuated. **Wear goggles at all times. The evacuated glassware can implode!**
6. Fill the boiling flask about 1/3 full with distilled water and fresh boiling chips. This will be used for the entire experiment.

7. When the thermometer is correctly positioned, the bottom of the bulb is even with the top of the boiling flask.
8. The condenser must be run during the experiment.
9. With the water aspirator, evacuate the system until the lowest pressure is reached (~ 60 - 80 torr). Then, close the stopcock on the trap.
10. Check for leaks by observing the manometer for 5 minutes after reaching the lowest pressure and closing the system. The system is leaking if the pressure changes by more than 1 or 2 torr in the five minute period. If the system is leaking, increase the pressure by opening the stopcock in the connection between the ballast and the trap. Then check all joints. The joints should be clean, greased and tight. Recheck the system for leaks by evacuating the system to lowest pressure again.
11. When the system is free of leaks, begin the heating process. Set the variastat at ~ 60 - 80 and heat rapidly until boiling. Then lower the variastat to ~ 30 to attain a gentle but steady boil.
12. Begin with the ascending cycle first. Take 10 readings spaced approximately as shown in the book. Then descend taking another 10 readings. Repeat this process one additional time. The aspirator should be off during the ascending cycle.
13. At each pressure allow the system to equilibrate. Take pressure/temperature readings simultaneously every 30 seconds. When three readings in a row are the same, the system has reached equilibrium.
14. After the completion of the experiment, release the vacuum, shut off the water to the condenser, empty the trap and wash the boiling flask. Do not put boiling chips in the drain. Dispose of the boiling chips in the trash.

16.3.3 Laboratory Report

1. Error Analysis:
 - (a) Estimate the error bars for only the first and last points of the graph used to determine the enthalpy of vaporization.
 - (b) Estimate the error in the slope using the limiting slope method.
 - (c) Estimate the error in Z by visual inspection of Figure 1 in the textbook.
 - (d) Propagate the errors to obtain an error in ΔH_v .
2. Conclusions: Compare your result with literature values. Use your experimentally derived enthalpy of vaporization to calculate the entropy of vaporization. Compare your result with Trouton's Rule [Engel and Reid, p. 194] and explain the differences.

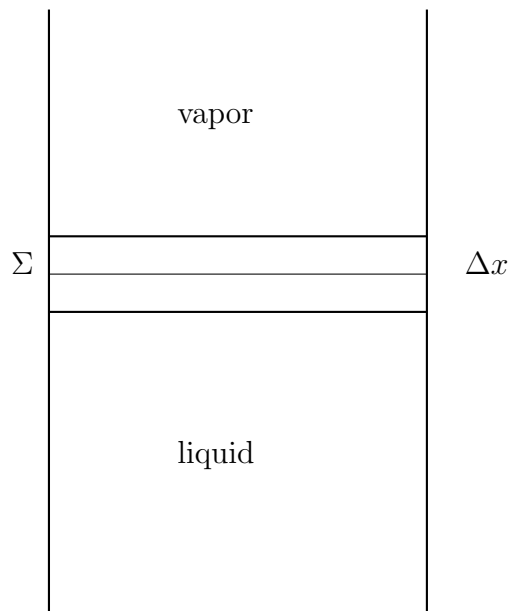


Figure 13: A representation of a container with a liquid mixture in equilibrium with a vapor phase. The section labelled Δx represents the surface region where the concentrations of both solvent and solute are not uniform. The horizontal line drawn at Σ represents the plane surface that defines the two-dimensional surface associated with the surface properties of the system.

16.4 Experiment 8(25) - Surface Tension of Solutions

16.4.1 Reading Assignment

Textbook, Experiment number 8(25), pages 95-104.

16.4.2 Theory

1. Description of the system: We consider a system in a container consisting of a solute dissolved in a liquid solvent so that the solute has bulk concentration c_s , and we also consider the situation where the solution has a surface of area σ . We take the surface to be in contact with a vapor phase (see Fig. 13). We assume the external temperature T and pressure P are fixed. While the solute has uniform density throughout most of the liquid phase, there is a small volume in the vicinity of the surface (represented by Δx in Fig. 13) where the solute and solvent concentrations decay to their vapor-phase concentrations. The concentration in the vapor phase is sufficiently low compared to the liquid phase that we can neglect further consideration of the contribution from the vapor phase to the surface properties. Figure 14 is a representation of the solute and solvent concentrations as a function of distance roughly within Δx . The x axis of Fig. 14 represents the vertical distance such that negative coordinates that are large in magnitude are locations within the bulk liquid, and large positive values of x (i.e. > 10) are in the vapor phase. The actual numerical values of x in the figure have no real meaning. The y axis of Fig. 14 gives the concentrations in arbitrary units for both the solvent and solute as marked. In Fig. 14 the bulk solute concentration is taken to be half of the solvent concentration, but the exact difference is unimportant for the current discussion. Of course we do require the solvent to have a greater bulk concentration than the solute. In the vicinity of the surface the solvent concentration decreases from its bulk value and the solute concentration increases. At $x = 10$ (arbitrary units) both

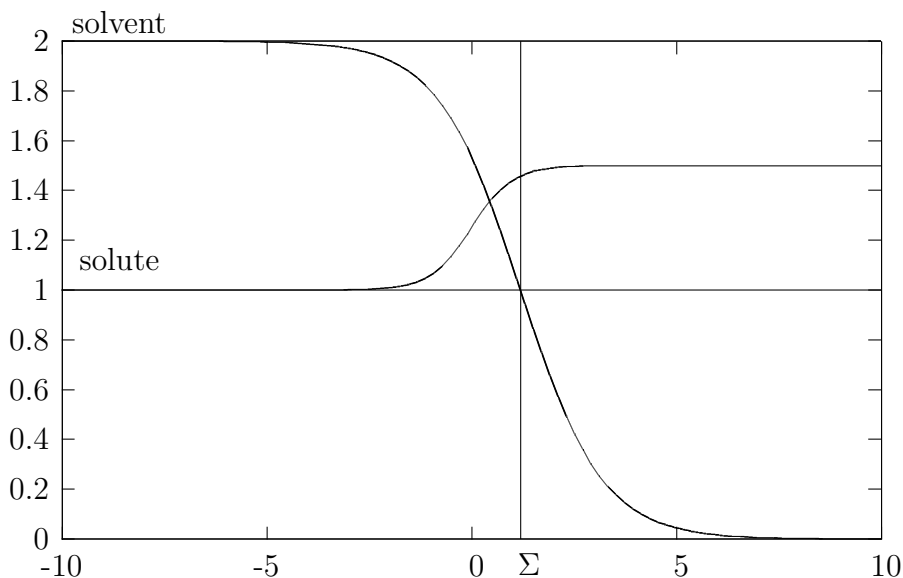


Figure 14: The concentration of the solvent and solute as a function of vertical distance (see Fig. 13) in the region Δx . The vertical line at Σ represents the constructed dividing surface. The horizontal lines at $y = 1$ and 2 (arbitrary units) correspond to the bulk solute and solvent concentrations respectively. The boundary at $x = 10$ (arbitrary units) is the end of the boundary region and corresponds to the interface between liquid and vapor.

the solute and solvent concentrations rapidly decay to their vapor phase values. In this experiment where we study a mixture of water (solvent) and butanol (solute), we have this behavior, which decreases the overall surface tension of the system. For other systems, like NaCl in water the solute concentration decreases near the surface and the solvent concentration increases.

2. Definition of the dividing surface: We next need to define a dividing surface that provides a sensible and useful meaning to the notion of the surface concentrations of solute and solvent in the region represented by Δx in Fig. 13. Shortly, we define the excess surface concentration of solute and solvent, and the units of the excess surface concentration are moles per unit area. The units imply that we can think of the solute and solvent molecules in Δx as residing on a two-dimensional planar surface, and we need to define our concepts so that the resulting definitions both give some physical insight and simultaneously are consistent with the number of moles of each species in the volume defined by Δx . The resulting definitions are physically artificial but mathematically consistent.

We draw a plane dividing surface within Δx , which we label Σ , and we choose the location of Σ so that the excess number of moles of solvent on each side of Σ is 0. To understand the notion of the excess number of surface moles of solvent, we let c_S be the bulk solvent concentration assumed to be uniform from the bottom of the container

to the constructed surface at Σ . We let \bar{c}_S be the actual surface concentration that is uniform from the bottom of the container until we reach Δx and decays in Δx as we approach the vapor phase. The behavior of \bar{c}_S is depicted in Fig. 14. The excess number of moles of solvent is defined to be the difference between the deficit in the number of moles of solvent obtained from \bar{c}_S compared to c_S below and above Σ . To make the concept clearer, we define the volume of the liquid phase, V_ℓ to be the volume of solution from Σ to the bottom of the container. Then the excess number of moles of solvent, n_S^Σ is defined to be

$$n_S^\Sigma = c_S V_\ell - \sigma \int_{-b}^{\Sigma} \bar{c}_S dx - \sigma \int_{\Sigma}^{\infty} \bar{c}_S dx \quad (125)$$

where $-b$ is the coordinate at the bottom of the container. We choose the location of Σ so that $n_S^\Sigma = 0$. This choice of the location of Σ implies

$$c_S V_\ell = \sigma \int_{-b}^{\infty} \bar{c}_S dx; \quad (126)$$

i.e. the total number of moles of solvent obtained by integrating the actual concentration profile over the entire system is identical to the number of moles that would be obtained if a uniform concentration existed up to the surface defined by Σ .

We define the excess number of moles of solute in the same manner used to define the solvent. Letting c_s be the bulk solute concentration assumed to be uniform throughout the solution and \bar{c}_s be the actual solute concentration, we write for the excess number of moles of solute

$$n_s^\Sigma = c_s V_\ell - \sigma \int_{-b}^{\Sigma} \bar{c}_s dx - \sigma \int_{\Sigma}^{\infty} \bar{c}_s dx. \quad (127)$$

While we have defined Σ so that $n_S^\Sigma = 0$ by construction, in general the different concentration profile for the solute implies that n_s^Σ cannot generally vanish.

As discussed above, we now associate (i.e. pretend) that all the excess number of moles of solute n_s^Σ are confined to a two-dimensional layer at Σ . We then define the *excess surface concentrations* of both the solute and solvent by

$$\Gamma_s = \frac{n_s^\Sigma}{\sigma} \quad (128)$$

and

$$\Gamma_S = \frac{n_S^\Sigma}{\sigma}. \quad (129)$$

By construction $\Gamma_S = 0$, and our goal is to measure Γ_s assuming all the excess number of moles of solute reside on the surface Σ . The measurement is possible by using the *Gibbs isotherm*.

Before deriving the Gibbs isotherm it is useful to understand the notion of excess surface concentration with some definite physical models. We consider the case where

Δx consists only of a single layer of solute molecules adsorbed onto a solvent. Such films are called *monolayers*, and an example of such a system is oleic acid adsorbed on water. In such a case we place Σ just at the interface between the solvent and solute. Clearly, the excess surface concentration of solvent is 0, because the solvent molecules are completely absent from the monolayer. The resulting solute excess surface concentration is just the number of moles of solute in the monolayer per unit area. We see that our definition of excess surface concentration of solute is consistent with our expectations for a monolayer. In another limit we consider the case that solute and solvent are very similar (e.g. two completely miscible alcohols). In such a case Γ_s can be expected to be very small; i.e. very little solute needs to be associated with our artificial surface layer, because the concentration profiles of both solute and solvent in Δx have nearly the same functional form.

In the intermediate case between the two extremes outlined in the previous paragraph, the excess number of moles of solute can either be enhanced or diminished compared to $c_s V_\ell$. Work is required to form the surface owing to the breaking of bonds at the interface (e.g. the hydrogen bonds for liquid water), and solutes having weaker bonds than the solvent tend to have positive excess surface concentrations (e.g. butanol). Using similar reasoning solutes having stronger bonds than the solvent tend to have negative excess surface concentrations (e.g. sodium chloride).

3. The Gibbs Isotherm: We let G_v be the Gibbs free energy of the vapor phase, G_ℓ be the Gibbs free energy of the liquid phase and G_Σ be the Gibbs free energy of the surface phase. We allow all three phases to contain both solvent and solute. We can write for the total differential of the entire Gibbs free energy function

$$dG = dG_v + dG_\ell + dG_\Sigma. \quad (130)$$

In each phase G is a function of T, P and the number of moles of both the solvent and solute. Additionally, G_Σ depends on the surface area σ . We can then write (at fixed temperature and pressure)

$$dG = \mu_s dn_{v,s} + \mu_S dn_{v,S} + \mu_s dn_{\ell,s} + \mu_S dn_{\ell,S} + \mu_s dn_{\Sigma,s} + \mu_S dn_{\Sigma,S} + \gamma d\sigma. \quad (131)$$

where $n_{v,s}$ represents the number of moles of solute in the vapor, $n_{v,S}$ represents the number of moles of solvent in the vapor, $n_{\ell,s}$ represents the number of moles of solute in the liquid, $n_{\ell,S}$ represents the number of moles of solvent in the liquid, $n_{\Sigma,s}$ represents the number of moles of solute associated with the surface region, $n_{\Sigma,S}$ represents the number of moles of solvent associated with the surface region, and γ is the reversible work per unit area to form the surface; i.e. the surface tension. In the vapor and liquid phases alone, we know that

$$dG_v = \mu_s dn_{v,s} + \mu_S dn_{v,S} \quad (132)$$

and

$$dG_\ell = \mu_s dn_{\ell,s} + \mu_S dn_{\ell,S} \quad (133)$$

so that Eq. (131) becomes

$$dG_{\Sigma} = \mu_s dn_{\Sigma,s} + \mu_S dn_{\Sigma,S} + \gamma d\sigma. \quad (134)$$

Because the number of moles of each species and the surface area are all extensive variables with the chemical potentials and the surface tension intensive variables, by Euler's theorem on homogeneous functions we can write

$$G_{\Sigma} = \mu_s n_{\Sigma,s} + \mu_S n_{\Sigma,S} + \gamma \sigma. \quad (135)$$

After taking the total differential of Eq. (135) and subtracting Eq. (134), the result is a Gibbs-Duhem relation

$$-\sigma d\gamma = n_{\Sigma,s} d\mu_s + n_{\Sigma,S} d\mu_S. \quad (136)$$

or solving for the differential of the surface tension

$$-d\gamma = \Gamma_s d\mu_s + \Gamma_S d\mu_S \quad (137)$$

where we have recognized the excess surface concentrations as the number of moles of solute or solvent per unit area. As noted previously $\Gamma_S = 0$ by construction, and we know that

$$\mu_s = \mu^{\circ} + RT \ln a_s \quad (138)$$

where a_s is the activity of the solute defined so that the form of the chemical potential matches the ideal gas form. We can also write the activity in terms of the solute concentration

$$a_s = c_s \phi_s \quad (139)$$

where ϕ_s is called the activity coefficient. The activity coefficient can be shown to approach unity as the concentration of solute becomes low. Taking the differential of the chemical potential

$$d\mu_s = RT d \ln a_s \quad (140)$$

and substituting the differential into Eq. (137), we obtain

$$-d\gamma_s = RT \Gamma_s d \ln a \quad (141)$$

or

$$\Gamma_s = -\frac{1}{RT} \frac{d\gamma}{d \ln a_s}, \quad (142)$$

which is the Gibbs isotherm. The excess solute surface concentration Γ_s depends on the slope of the surface tension with solute activity, and to the extent that the relation between the surface tension and the logarithm of the activity is linear, we can assume the excess solute concentration is independent of total concentration. Consequently, the Gibbs isotherm can provide a single excess solute surface concentration for all total solution concentrations. The location of the dividing surface Σ does vary with concentration, but the location of Σ is not important to the overall interpretation.

In the textbook Eq. (142) is approximated by replacing the activity with the concentration; i.e. setting $\phi = 1$. Equation (142) enables the determination of the excess surface concentration of solute by measuring the surface tension as a function of concentration. We determine the excess surface concentration of solute in this manner in the current experiment.

16.4.3 Experimental Notes

1. We use the capillary rise method discussed in the textbook.
2. You make solutions of either *n*-butanol, iso-butanol, *sec*-butanol or *tert*-butanol. Ask your TA to assign the particular alcohol.
3. You are not to perform the NaCl experiment.
4. The approximately 0.8 M butanol stock solution will be provided. You are to use volumetric methods to prepare the other solutions.
5. Be sure to record the temperature of the bath.
6. The surface tension of distilled water as a function of temperature is given in the Table below. To determine the surface tension at your working temperature, you must use interpolation methods.

Temperature ($^{\circ}\text{C}$)	surface tension of water (dy cm^{-1})
10.	74.23
20.	72.75
25.	72.00
30.	71.20

16.4.4 Laboratory Report

The laboratory report format is as follows:

1. Title Page: Give the title of the experiment, and the date the experiment was performed.
2. Abstract: This should be a 1 paragraph summary of what is to follow including the results. It is suggested that you write the abstract after the rest of the report is complete.
3. Introduction: This should be a discussion of the purpose of the experiment.
4. Theory: This should provide a derivation of the key formulas used in the experiment and in the data analysis section. Definitions must be provided for all symbols used in this and subsequent sections.

5. Procedure: Attach the outline of the procedure you brought to class and had signed by your TA.
6. Original Data: Attach the original data signed by your TA.
7. Data Table: Tabulate your data so that it can be used in the calculations to follow.

16.4.5 Calculations

1. We assume that the density of water and butanol solutions are the same and unchanging during the course of the experiment. We can then determine the surface tension of each solution, γ_s using

$$\gamma_s = \gamma_{H_2O} \frac{h_s}{h_{H_2O}} \quad (143)$$

where γ_{H_2O} is the surface tension of water obtained from the table, h_s is the capillary rise of the solution and h_{H_2O} is the capillary rise of pure water.

2. Be sure to calculate the surface concentration in both units of mol cm⁻² and molecules per square angstrom as discussed in your textbooks.
3. For the butanol calculations, use the activity coefficients given on page 302 of your textbooks. Alternatively, you can use the graph of the activity coefficient as a function of concentration given Fig. 15. You may use the same activity coefficients for all forms of butanol.
4. Error Analysis: Propagate the errors for only the first and last points on the graph of the surface tension as a function of the logarithm of the activity of the solution. Use the two calculated errors to determine the limiting slopes of your graph. Use the limiting slopes to estimate the error in your determined surface concentrations. Make sure to express the errors in both mol cm⁻² and molecules Å⁻². A conversion is required.

17 Final Laboratory Assignment - Experiment Number 5(9) - Partial Molar Volume

17.1 Reading Assignment

Experiment 5(9) - Partial Molar Volume - Pages 72-28, text.

17.2 Experimental Notes

1. We use the Weld-type pycnometer. Read the note at the bottom of page 76, carefully.

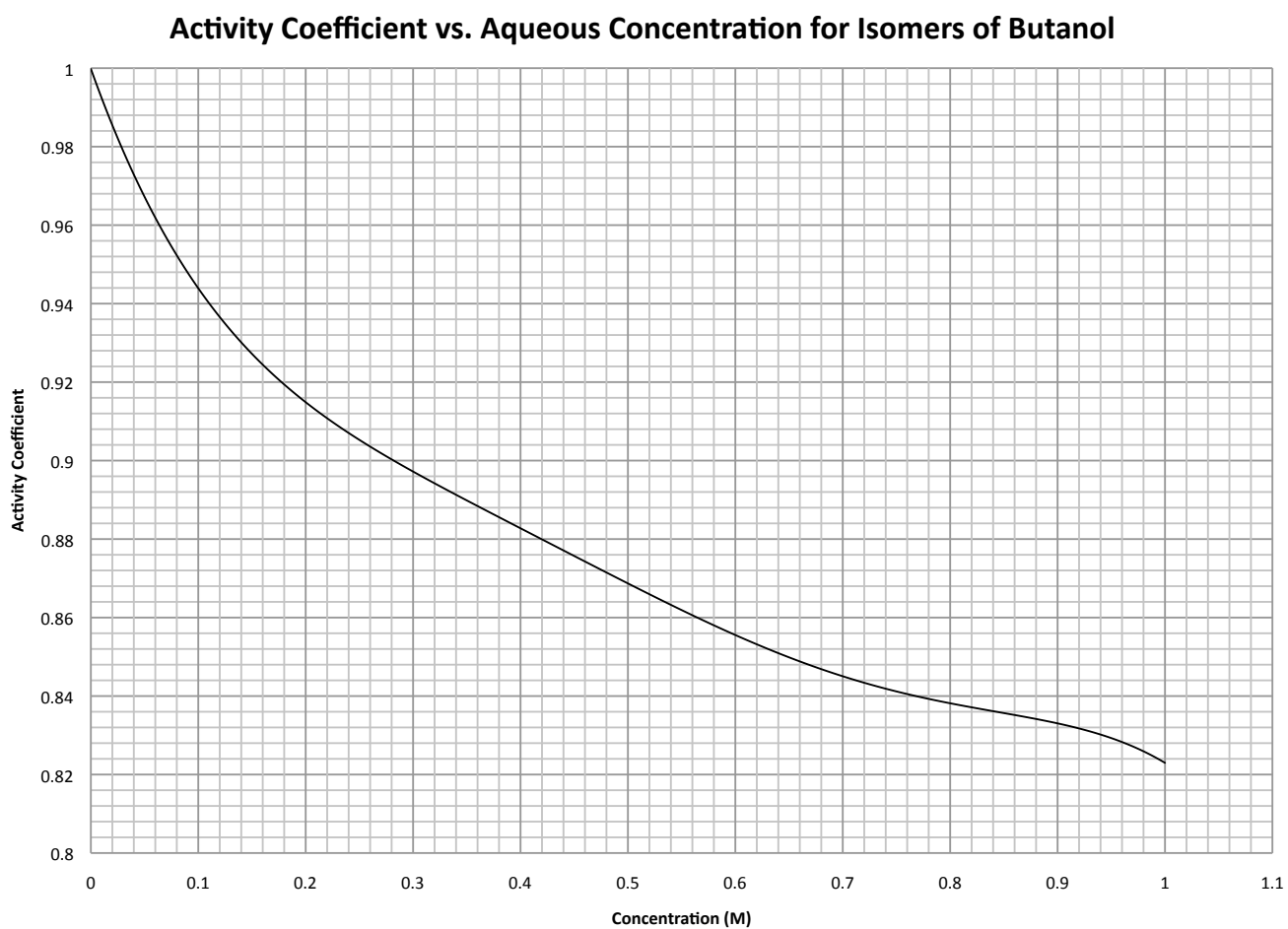


Figure 15: The activity coefficient of butanol as a function of concentration.

2. Weigh your empty pycnometer. Make sure it is dry and clean before the measurement and use the Mettler balance. Accurate measurement is essential. If you have any problems with the balance, ask your TA.
3. Fill the pycnometer with the distilled water contained in the bottle in the heat bath. Place the pycnometer in a beaker in the bath and wait five minutes to make sure the temperature is equilibrated. If necessary, add water to the pycnometer to be sure it is full.
4. Record the bath temperature and the room temperature. The temperature of the bath must be greater than the room temperature. If this is not true, inform your TA at once.
5. Remove the pycnometer from the bath, dry the exterior with a paper towel, and accurately weigh the full pycnometer with the Mettler balance.
6. Empty the pycnometer and dry it with acetone and the aspirator under the fume hood.
7. Fill the dry pycnometer with the stock NaCl solution in the flask held in the temperature bath. Place the pycnometer in a beaker in the bath for at least five minutes to equilibrate the temperature. Again add stock solution to make sure the pycnometer is full. Be sure to record the molality of the stock solution and its associated error.
8. Remove the pycnometer, dry the exterior and weigh the pycnometer carefully.
9. Clean and dry the pycnometer with water and acetone as above.
10. Pipette 25.0 ml of the stock solution into a 50.0 ml volumetric flask. Add distilled water from the heat bath slowly while mixing. Do not fill completely to the graduation mark. Place the flask in the heat bath to equilibrate for 5 minutes. After equilibration, add distilled water from the bottle in the heat bath to the graduation mark.
11. Place this solution in the pycnometer and equilibrate using the same procedure outlined above.
12. After equilibration, weigh the pycnometer with diluted solution carefully as above. After the determination of the weight, do not dispose the diluted solution. Return the solution to the volumetric flask. You will need the solution for the next dilution.
13. Repeat the process for a further dilution of the diluted NaCl solution by an additional 50%.
14. Clean all glass ware before leaving the laboratory.

17.3 Laboratory Reports

Your reports are to be written in class during Week 13. Come to 472 Beaupre for the report writing session. You should bring calculators, your original signed data, your original procedure, signed by your TA and pens or pencils. We supply paper and graph paper. We also supply a copy of pages 72-78 from your text books as well as a copy of these notes. You may not use your own texts. **We strongly recommend that you prepare your laboratory reports at home for practice. Although you are not allowed to bring your home written reports with you to the report writing session, the practice is essential in order to be able to finish the report in the allotted time (4 hours).** You may come to any one of the instructors in the course for help prior to the report writing session. In fact, we encourage you to come to us for help. Success in this assignment will only be possible with thorough preparation.

The laboratory report format is as follows:

1. Title Page: Give the title of the experiment, and the date the experiment was performed.
2. Abstract: This should be a 1 paragraph summary of what is to follow including the results. It is suggested that you write the abstract after the rest of the report is complete.
3. Introduction: This should be a discussion of the purpose of the experiment.
4. Theory: This should provide a derivation of the key formulas used in the experiment and in the data analysis section. Definitions must be provided for all symbols used in this and subsequent sections.
5. Procedure: Attach the outline of the procedure you brought to class and had signed by your TA.
6. Original Data: Attach the original data signed by your TA.
7. Results: Tabulate your data so that it can be used in the calculations to follow.
8. Calculations: Perform all calculations to determine ϕ and \sqrt{m} for the three measured concentrations. Please note that the relation between the molarity M of the solution and the molality m is given by Eq. (24) in your textbook. Because the units associated with Eq. (24) of the textbook are confusing, a derivation of the connection now follows.

Consider a solution containing 1000. grams of solvent of known molality m . Following the notation in the textbook, we let M_2 represent the molecular weight of the solute in units of grams per mole. Then the total mass in grams of solute and solvent in the solution of molality m is $1000. \text{ g} + mM_2$. The density d of the solution is then given by

$$d = \frac{1000. \text{ g} + mM_2}{V} \quad (144)$$

where V is the volume of the solution in cm^3 . By definition the molarity M is the number of moles of solute in 1 liter of solution, so that the volume of the solution in liters is given by

$$\Omega = \frac{\text{L}}{1000. \text{ cm}^3} \frac{1000. \text{ g} + mM_2}{d}. \quad (145)$$

Because the number of moles of solute in our solution of 1000. grams of solvent is m , the molarity is given by

$$M = \frac{m}{\Omega} \quad (146)$$

$$= \frac{1000. \text{ cm}^3}{\text{L}} \frac{md}{1000. \text{ g} + mM_2} \quad (147)$$

or

$$M = \left(\frac{1000. \text{ cm}^3}{\text{L}} \right) d \left(\frac{1}{\frac{1000. \text{ g}}{m} + M_2} \right). \quad (148)$$

If we solve Eq.(148) for m in terms of M , we obtain

$$m = 1000. \text{ g} \frac{1}{\left(\frac{d}{M} \right) \left(\frac{1000. \text{ cm}^3}{\text{L}} \right) - M_2}. \quad (149)$$

Note that in the experiment M is decreased by a factor of 2 for each dilution. The relation for m from the dilutions is more complex and must be calculated using Eq. (149).

9. Graphical Analysis: Plot your determined values of ϕ as a function of \sqrt{m} . You should find a linear relationship. Plot the best straight line through the data points and determine the slope and the intercept ϕ^0 . Please note that the book uses sloppy notation for the slope. In Eqs. (21) and (22) you should interpret $d\phi/d\sqrt{m}$ as the derivative (i.e. slope) of your plotted line. The d represents differentiation in this case. In Eqs. (17) and (24) the d is the density.
10. Error Analysis: Propagate errors to determine the errors in ϕ and \sqrt{m} for the lowest and the highest concentrations. Show all calculations. Tabulate all final results with associated errors. Plot the error bars on your graph to determine the errors in the slope and intercept using the method of limiting slopes.

Hint: The formulas from which you determine error may appear complex. However, if you break the problem of determining errors into parts, the task can be made relatively easy. By breaking the problem into parts we introduce a small error that can be ignored for the purposes of this assignment. As an example, consider Eq. (17) in your text

$$\phi = \frac{1}{d} \left(M_2 - \frac{1000}{m} \frac{W - W_0}{W_0 - W_e} \right).$$

For our purposes, we can assume the error in M_2 is too small to affect our results. Take the value of M_2 to be $58.443 \text{ g mol}^{-1}$ with no error. Define

$$a = W - W_0$$

and

$$b = W_0 - W_e$$

The errors in a and b are given by

$$\epsilon(a) = \sqrt{\epsilon(W)^2 + \epsilon(W)_0^2}$$

and

$$\epsilon(b) = \sqrt{\epsilon(W_0)^2 + \epsilon(W_e)^2}$$

Now let

$$c = \frac{1000}{m} \frac{a}{b}$$

The error in c is given by

$$\frac{\epsilon(c)}{c} = \sqrt{\left(\frac{\epsilon(a)}{a}\right)^2 + \left(\frac{\epsilon(b)}{b}\right)^2 + \left(\frac{\epsilon(m)}{m}\right)^2}$$

Finally, the error in ϕ is given by

$$\frac{\epsilon(\phi)}{\phi} = \sqrt{\left(\frac{\epsilon(c)}{M_2 - c}\right)^2 + \left(\frac{\epsilon(d)}{d}\right)^2}$$

For the other equations, use the same methods to break the problem of calculating errors into parts.

11. Summary of Final Results: Report the values of the slope and intercept with the associated errors. Be sure to report the proper number of significant figures and the units. Use Eqs. (21), (22), the determined values of ϕ^0 and the slope to calculate the partial molar volumes of water and NaCl as a function of m for $m=0.5, 1.0, 1.5, 2.0$ and 2.5 . Plot the results for the partial molar volumes on a common graph.
12. Conclusions: Summarize your results.

18 Tables of Data

ORGANIC LIQUIDS*
 N. ERNEST DORSEY

Viscosity of Water

FORMULAE AND UNITS

At a pressure of 1 atm., $\eta = a/(b + t)^2$.

At a pressure of P kg/cm², $\eta_p = \eta(1 + k_c(P - 1) \times 10^{-4})$.

η_1 is the value of η when P is 1 kg/cm², which may be taken as the value of η at 1 atm.

The unit of η is the poise unless otherwise stated.

WATER BETWEEN 0 AND 100°. I. C. T. VALUES

The following table was prepared from a critical evaluation of all available data. It is estimated that the accuracy is of the order of 0.1% between 0 and 40° and of 0.5 to 1% at higher temperatures. Linear interpolation may be safely employed throughout the table.

Values in millipoises (1, 12, 16, 17, 22, 24, 30, 31, 32, 38)

$t, ^\circ\text{C}$	0	1	2	3	4	5	6	7	8	9
0	17.938	17.320	16.740	16.193	15.676	15.188	14.726	14.288	13.872	13.476
10	13.097	12.735	12.390	12.061	11.748	11.447	11.156	10.875	10.603	10.340
20	10.087	9.843	9.608	9.380	9.161	8.949	8.746	8.551	8.363	8.181
30	8.004	7.834	7.670	7.511	7.357	7.208	7.064	6.925	6.791	6.661
40	6.536	6.415	6.298	6.184	6.075	5.970	5.868	5.770	5.675	5.582
50	5.492	5.405	5.320	5.236	5.153	5.072	4.994	4.918	4.843	4.770
60	4.699	4.629	4.561	4.495	4.431	4.368	4.306	4.245	4.186	4.128
70	4.071	4.016	3.962	3.909	3.857	3.806	3.756	3.708	3.661	3.615
80	3.570	3.526	3.483	3.440	3.396	3.357	3.317	3.278	3.240	3.203
90	3.166	3.130	3.095	3.061	3.027	2.994	2.962	2.930	2.899	2.869
100	2.839	2.82	2.79	2.76	2.73	2.70	2.67	2.64	2.62	2.59

ABSOLUTE DENSITY OF WATER

DENSITY IN GRAMS PER CUBIC CENTIMETER, COMPUTED FROM THE RELATIVE VALUES BY THIESEN, SCHEEL AND DISSELHORST (1900), AND THE ABSOLUTE VALUE AT 3.98° C. BY THE INTERNATIONAL BUREAU OF WEIGHTS AND MEASURES (1910).

Degrees	0	1	2	3	4	5	6	7	8	9
0	0.999841	847	854	860	866	872	878	884	889	895
1	900	905	909	914	918	923	927	930	934	938
2	941	944	947	950	953	955	958	960	962	964
3	965	967	968	969	970	971	972	972	973	973
4	973	973	973	972	972	972	970	969	968	966
5	965	963	961	959	957	955	952	950	947	944
6	941	938	935	931	927	924	920	916	911	907
7	902	898	893	888	883	877	872	866	861	855
8	849	843	837	830	824	817	810	803	796	789
9	781	774	766	758	751	742	734	726	717	709
10	700	691	682	673	664	654	645	635	625	615
11	605	595	585	574	564	553	542	531	520	509
12	498	486	475	463	451	439	427	415	402	390
13	377	364	352	339	326	312	299	285	272	258
14	244	230	216	202	188	173	159	144	129	114
15	099	084	069	054	038	023	007	*991	*975	*959
16	0.998943	926	910	893	877	860	843	826	809	792
17	774	757	739	722	704	686	668	650	632	613
18	595	576	558	539	520	501	482	463	444	424
19	405	385	365	345	325	305	285	265	244	224
20	203	183	162	141	120	099	078	056	035	013
21	0.997992	970	948	926	904	882	860	837	815	792
22	770	747	724	701	678	655	632	608	585	561
23	538	514	490	466	442	418	394	369	345	320
24	296	271	246	221	196	171	146	120	095	069
25	044	018	*992	*967	*941	*914	*888	*862	*836	*809
26	0.996783	756	729	703	676	649	621	594	567	540
27	512	485	457	429	401	373	345	317	289	261
28	232	204	175	147	118	089	060	031	002	*973
29	0.995944	914	885	855	826	796	766	736	706	676
30	646	616	586	555	525	494	464	433	402	371

A Appendix: The Critical Exponent of the Van der Waals Fluid Along the Critical Isobar

To derive the critical exponent of C_p along the critical isobar we start with Eq. (96) and substitute the pressure and temperature and volume using the reduced variables

$$P_r = \frac{P}{P_c} \quad (150)$$

$$T_r = \frac{T}{T_c} \quad (151)$$

and

$$V_r = \frac{V}{V_c} \quad (152)$$

resulting in the reduced equation of state

$$P_r = \frac{8T_r}{3V_r - 1} - \frac{3}{V_r^2}. \quad (153)$$

When the Van der Waals equation of state is expressed in reduced variables as in Eq. (153), the result is independent of a and b ; i.e. the equation of state is the same for all gases. This universal behavior is often called the *Principle of Corresponding States*. We next define

$$\tau = T_r - 1 \quad (154)$$

$$\nu = V_r - 1 \quad (155)$$

and along the critical isobar we set $P_r = 1$, so that we have defined the critical isobar. The resulting relation between τ and ν is

$$\tau = \frac{1}{8} \left(\frac{3\nu^3 + 7\nu^2 + 14\nu + 7}{(\nu + 1)^2} \right). \quad (156)$$

At the critical point both τ and ν vanish.

To determine how ν depends on τ as the critical point is approached along the critical isobar, we differentiate Eq. (156) to obtain

$$\left(\frac{\partial \tau}{\partial \nu} \right)_{P_r=1} = \frac{3}{8} \frac{\nu^3}{(\nu + 1)^3} + \frac{9}{8} \frac{\nu^2}{(1 + \nu)^3} \quad (157)$$

Then

$$\lim_{\tau \rightarrow 0} \left(\frac{\partial \tau}{\partial \nu} \right)_{P_r=1} = \frac{3}{8} \nu^3 + \frac{9}{8} \nu^2 \quad (158)$$

We next integrate both sides of Eq.(158) to obtain

$$\tau = \frac{3}{32} \nu^4 + \frac{3}{8} \nu^3 \quad (159)$$

and

$$\lim_{\nu \rightarrow 0} \tau = \frac{3}{8} \nu^3 \quad (160)$$

or

$$\lim_{\tau \rightarrow 0} \nu = \frac{2}{3^{1/3}} \tau^{1/3}. \quad (161)$$

We next express Eq.(117) in terms of the critical temperature, pressure and volume using

$$T = T_r T_c \quad V = V_r V_c \quad P = P_r P_c$$

to give

$$\frac{C_P - C_V}{R} = \frac{P_c V_c T_r}{R T_c} \left(\frac{\partial V_r}{\partial T_r} \right)_{P_r} \left(\frac{\partial P_r}{\partial T_r} \right)_{V_r} \quad (162)$$

$$= \frac{3}{8} T_r \left(\frac{\partial V_r}{\partial T_r} \right)_{P_r} \left(\frac{\partial P_r}{\partial T_r} \right)_{V_r}. \quad (163)$$

From the reduced equation of state

$$\left(\frac{\partial P_r}{\partial T_r} \right)_{V_r} = 4 \quad (164)$$

and using Eqs. (154), (155) and (161)

$$\lim_{\tau \rightarrow 0} \frac{C_P - C_V}{R} = \lim_{\tau \rightarrow 0} \frac{3}{2} (\tau + 1) \left(\frac{\partial \nu}{\partial \tau} \right)_{P_r=1} \quad (165)$$

$$= \frac{1}{3^{1/3}} \tau^{-2/3} \quad (166)$$

demonstrating that along the critical isobar the C_P critical exponent is $2/3$.